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A graphical language for quantum protocols based on the category of cobordisms

Dušan ĐorĐević 💿 · Zoran Petrić · Mladen Zekić

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Abstract As shown by Abramsky and Coecke, quantum mechanics can be studied in terms of dagger compact closed categories with biproducts. Within this structure, many well-known quantum protocols can be described and their validity can be shown by establishing the commutativity of certain diagrams in that category. In this paper, we propose an explicit realization of a category with enough structure to check the validity of a certain class of quantum protocols. To do this, we construct a category based on one-dimensional cobordisms with attached elements of a certain group freely generated by a finite set. We use this category as a graphical language, and we show that it is dagger compact closed with biproducts. Then relying on the coherence result for compact closed categories, proved by Kelly and Laplaza, we show the coherence result, which enables us to check the validity of quantum protocols just by drawing diagrams. In particular, we show the validity of quantum teleportation, entanglement swapping (as formulated in the work of Abramsky and Coecke) and superdense coding protocol.

Keywords Dagger compact closed category · Biproducts · Cobordism · Entanglement · Coherence

Mathematics Subject Classification 18M10 · 18M40 · 18D20 · 81P45 · 57Q20

1 Introduction

This paper offers a mathematical result and its application in the field of quantum information. The goal is to provide a minimal graphical language sufficient for the verification of categorical quantum protocols. This is achieved through a category $1\mathbf{Cob}_{\mathfrak{G}}^{\oplus}$ based on the category of one-dimensional cobordisms. We will work out in some detail a few quantum protocols and apply a technique of their verification based on the category $1\mathbf{Cob}_{\mathfrak{G}}^{\oplus}$. This work contains both pure mathematical results as well as results in the applied field of quantum information, which are related to

Z. Petrić · M. Zekić Mathematical Institute SANU, Knez Mihailova 36, p.f. 367, 11001 Belgrade, Serbia e-mail: zpetric@mi.sanu.ac.rs

M. Zekić e-mail: mzekic@mi.sanu.ac.rs

D. ĐơrĐević (⊠) Faculty of Physics, Studentski trg 12, 11001 Belgrade, Serbia e-mail: dusan.djordjevic@ff.bg.ac.rs

both physics and computer science. Therefore, we will try to make our exposition sufficiently detailed to make the results available to a larger group of scientists from different areas.

There are several diagrammatic calculi proposed earlier (e.g., [7,31]) appropriate for the verification of categorical quantum protocols. Our intention is to make one within a frame of a category relevant to quantum mechanics—the category of cobordisms. This is achieved by relying on a result from Ref. [20] just by relaxing (i.e., neglecting some components of) the arrows of a category constructed in that paper. However, some labels of connected components of the underlying manifolds remain in our calculus, but we have minimized their role. One way to get rid of all labels is to increase the dimension of cobordisms in question. We discuss this suggestion in Sect. 9. Note that we are not trying to provide a diagrammatic calculus that is better for practical purposes than the existing ones (for example, ZX calculus [8]), but rather to, in a very precise manner, formulate one calculus that uses geometry and that is able to capture some aspects of quantum mechanics. We are aware that our method is not universal and that the complexity of certain quantum phenomena cannot be captured by one-dimensional cobordisms. Also, we do not claim that we can formulate and check the validity of all quantum protocols (for example, it is hard to simulate the most general unitary operator acting on two qubits). A language based on cobordisms should be taken as a mathematical tool having some computational advantages compared to standard "linear-word" languages. However, as usual with syntax, such a language does not capture all the mathematical details of the subject.

Ever since its formulation in the first half of the twentieth century, quantum mechanics is naturally set to live in a separable Hilbert space. This enables one to talk about notions such as entanglement and measurement in an almost trivial way. However, they are far from being understood by the scientific community. While the basic mathematical formalism is easy to understand, its physical meaning is much less clear and various interpretations of quantum mechanics are possible, without currently any basis on which we could select only one that is correct. This means that it could be fruitful to reconsider some basic notions about quantum mechanics and to try to formulate it in terms of a different mathematical structure. For example, a recently emerged and fast expanding field of *fractional quantum mechanics* is developed by applying the *fractional calculus* (which is a generalization of classical differential and integral calculus) in quantum physics. Following this approach, and in a similar spirit to our research, authors in the recent study, Ref. [3] simulated the spatial form of the fractional Schrödinger equation for the electrical screening potential using the fractional derivatives and the numerical simulation methods.

Similar motivation led authors of Ref. [1] to develop the so-called categorical approach to quantum mechanics in terms of dagger compact closed categories with biproducts. Of course, the starting point is naturally the case where standard Hilbert-space formalism leads to the consideration of finite dimensional spaces (for example, consideration of spin 1/2 gives rise to the two dimensional Hilbert space; on the other hand, it is not too hard to construct classical theories whose Hilbert space after quantization turns out to be finite dimensional, [33], although they are much less known for an average physicist). For simplicity, and for the practical application, we restrict to the case of two dimensional Hilbert spaces (qubits). Basis vectors are $|0\rangle$ and $|1\rangle$ (we use Dirac notation for this part). If we denote $\mathcal{H} = \text{span } (|0\rangle, |1\rangle$), then the total Hilbert space of a composite bipartite system is $\mathcal{H} \otimes \mathcal{H}$. As we use tensor product, there are no projections to \mathcal{H} , and we can introduce the notion of an entangled state. For example, we can form a basis in $\mathcal{H} \otimes \mathcal{H}$ from the entangled states as

$$\begin{split} |\beta_1\rangle &= \frac{1}{\sqrt{2}}(|0\rangle \otimes |0\rangle + |1\rangle \otimes |1\rangle), \\ |\beta_2\rangle &= \frac{1}{\sqrt{2}}(|0\rangle \otimes |1\rangle + |1\rangle \otimes |0\rangle), \\ |\beta_3\rangle &= \frac{1}{\sqrt{2}}(|0\rangle \otimes |1\rangle - |1\rangle \otimes |0\rangle), \\ |\beta_4\rangle &= \frac{1}{\sqrt{2}}(|0\rangle \otimes |0\rangle - |1\rangle \otimes |1\rangle). \end{split}$$

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Basis $\{\beta_i\}_{i=1}^{i=4}$ is usually referred to as Bell basis (see e.g., Ref. [27]). Of course, this construction is made on a few assumptions. The first one is a linear structure of vector spaces. Historically, linear structure was a natural guess based on an intuition about wavelike properties of particles (electrons). For example, it was known long before the birth of quantum mechanics that light can be described in terms of oscillating electric and magnetic fields and that for those fields, the superposition principle holds (linearity of Maxwell's equations). Moreover, the intensity of a wave is proportional to $|\mathbf{E}|^2$ (where **E** stands for a complex representative of the electric field), and this further motivates the Born rule for probabilities associated with the measurement outcome. Despite its success, it is still intriguing to consider theories without vector space structure.

Passing from Hilbert spaces to categorical semantics of quantum protocols means abstracting the superfluous structure and keeping only the properties of the category of Hilbert spaces necessary to express these protocols. There are no other assumptions about the categories which could cause suspicion during the process of protocol validation. This means that checking the validity of such protocols using categorical semantics could be instantiated in any type of categories satisfying these properties. Needless to say that it includes the category of Hilbert spaces, where quantum mechanics is formulated.

Quantum mechanics can be considered as a special case of quantum field theories for 0 + 1 dimensions. On the other hand, it is well known that one approach to quantum field theories (especially to the case of topological quantum field theories) is using cobordisms to represent space-time evolution processes. This opens another natural question, and that is to what extent one can use the category 1**Cob** (of one-dimensional cobordisms) to simulate quantum mechanical processes. In addition, there are ideas from quantum gravity, in the context of AdS/CFT correspondence [24], that some aspects of quantum theory (for example, entanglement entropy) could be obtained from geometry [30]. Quantum–geometry relation is also evident in ER = EPR proposal [25]. Therefore, seeking the role of geometry in quantum physics is interesting on its own. The results of our paper can be viewed as a step toward understanding the connection between quantum mechanics and geometry. We have to note that our ambitions are not to define the most useful computing tool for checking quantum protocols, but rather to show that there are cases where geometry and quantum physics combine, and where their relations can be formulated using very precise mathematics. See also Ref. [6] for some other relations between geometry and quantum physics.

To obtain this correspondence, we will introduce the notion of \mathfrak{G} -cobordisms. They correspond to regular cobordisms (for a precise definition of 1-dimensional cobordisms, see the following sections), but with additional structure, such that each connected component has an element of a group \mathfrak{G} attached. This introduces a notion of a \mathfrak{G} -segment or a \mathfrak{G} -circle. Group elements will play the role of (unitary) transformations that can be done on a quantum state. Note that a similar idea was discussed in Ref. [19], but without explicitly referring to the categorical quantum mechanics. In our applications to quantum protocols, we will use specific group elements that are suitable for applications in those protocols we describe, but one could, in general, use, without changing the main conclusions, other examples to describe other protocols.

On the other hand, a motivation for our work can be purely mathematical. In category theory and its application, it is of great importance to establish whether a certain diagram commutes. Usually, this is done by inspection, using a set of equalities (for example, as those from Appendix A). Though, in principle, a straightforward task, it usually consumes a non-negligible amount of time. For this reason, it is practical to prove certain coherence results. Such a result enables one to check the commutativity of diagrams, consisting of canonical arrows of a certain categorical structure, just by drawing pictures in an appropriate graphical language. It is clear that such a calculation by drawing pictures some effort, and our intention was to prepare such a graphical language that minimizes the possibility of errors and decreases the computational complexity.

A detailed explanation of our approach to coherence is given in [29, Introduction], where also results akin to those proven here are presented. Briefly, we start with a freely generated category built out of syntax material, whose objects are formulae and arrows are equivalence classes of terms in an equational system. Then we show a completeness result with respect to a model in the form of a graphical category.

In Sect. 2, we review some basic categorical notions relevant for this paper. In Sect. 3, we further discuss the category 1**Cob**. Section 4 introduces two compact closed categories with some additional structure both freely generated by a free group considered as a category. The isomorphism of these two categories is established in that

section. The next two sections (Sects. 5 and 6) are technical necessities and could be skipped in the first reading. In Sect. 7, we introduce gradually (in several steps) the category $1\mathbf{Cob}_{\mathfrak{G}}^{\oplus}$, and formulate the coherence result essential for the verification of certain categorical quantum protocols, a task we perform in Sect. 8. A possibility of omitting the labels via increasing the dimension of cobordisms is given in Sect. 9. We give our final remarks in Sect. 10. Appendix A contains an equational presentation of dagger compact closed categories with dagger biproducts and Appendix B discusses a categorical approach to scalars and probability amplitudes.

To conclude, this is a theoretical study. It is experimentally validated as far as the mentioned quantum protocols are. We expect more real-world examples, implementations, and applications of this theory after making an extension with three-dimensional cobordisms, which is an ongoing project.

2 Closed categories and biproducts

Some notions from category theory relevant for this paper are introduced in this section. A symmetric monoidal category is a category \mathcal{A} equipped with a distinguished object I, a bifunctor $\otimes : \mathcal{A} \times \mathcal{A} \to \mathcal{A}$ (we abbreviate $\mathbf{1}_a \otimes f$ and $f \otimes \mathbf{1}_a$ by $a \otimes f$ and $f \otimes a$, respectively) and the natural isomorphisms α , λ and σ with components $\alpha_{a,b,c}: a \otimes (b \otimes c) \to (a \otimes b) \otimes c$, $\lambda_a: I \otimes a \to a$ and $\sigma_{a,b}: a \otimes b \to b \otimes a$. (Note that in Ref. [1], λ denotes the inverse of our λ and due to the presence of symmetry, we do not introduce a name for the isomorphism $a \cong a \otimes I$.) Moreover, the coherence conditions concerning the arrows of \mathcal{A} (see the equalities A.6–A.8 in Appendix A) hold. A symmetric monoidal category is *monoidally strict* when the operation \otimes on its objects is associative with I being the neutral, and moreover, the arrows α and λ are identities.

A compact closed category is a symmetric monoidal category in which every object *a* has its dual a^* . This means that there are units $\eta_a : I \to a^* \otimes a$ and counits $\varepsilon_a : a \otimes a^* \to I$ such that the equalities A.9 of Appendix A hold. If a functor between two compact closed categories preserves this structure "on the nose", then we say that it strictly preserves the compact closed structure, and we use the same terminology in other cases.

It is straightforward to conclude that the following isomorphisms hold in every compact closed category.

$$u_{a,b}: (a \otimes b)^* \cong b^* \otimes a^*, \quad v: I^* \cong I, \quad w_a: a^{**} \cong a$$

(In a monoidally strict compact closed category,

$$u_{a,b} = (b^* \otimes a^* \otimes \varepsilon_{a \otimes b}) \circ (b^* \otimes \eta_a \otimes b \otimes (a \otimes b)^*) \circ (\eta_b \otimes (a \otimes b)^*)$$

 $v = \varepsilon_I$ and $w_a = (\varepsilon_{a^*} \otimes a) \circ (\sigma_{a^{**}, a^*} \otimes a) \circ (a^{**} \otimes \eta_a)$.) A compact closed category is *strict* when it is monoidally strict and $(a \otimes b)^* = b^* \otimes a^*$, $I^* = I$ and $a^{**} = a$, while $u_{a,b}$, v and w_a are identities.

For quantum protocols discussed below, the following derived operations on arrows of a compact closed category are frequently used. For $f: a \to b$, its name $\lceil f \rceil: I \to a^* \otimes b$ and its coname $\lfloor f \rfloor: a \otimes b^* \to I$ are defined as

$$\ulcorner f \urcorner = (a^* \otimes f) \circ \eta_a, \quad \llcorner f \lrcorner = \varepsilon_b \circ (f \otimes b^*).$$

The function * on objects of a compact closed category \mathcal{A} , extends to a functor $*: \mathcal{A}^{op} \to \mathcal{A}$ in the following way. For $f: a \to b$, let $f^*: b^* \to a^*$ be

$$\lambda_{a^*} \circ \sigma_{a^*,I} \circ (a^* \otimes \varepsilon_b) \circ \alpha_{a^*,b,b^*}^{-1} \circ ((a^* \otimes f) \otimes b^*) \circ (\eta_a \otimes b^*) \circ \lambda_{b^*}^{-1}.$$

A *dagger category* is a category \mathcal{A} equipped with a functor $\dagger: \mathcal{A} \to \mathcal{A}^{op}$ such that for every object *a* and every arrow *f* of this category, $a^{\dagger} = a$, and $f^{\dagger\dagger} = f$. (For more details, see Refs. [15,31].) A *dagger* compact closed category is a compact closed category \mathcal{A} , which is also a dagger category satisfying the equalities A.21–A.23 of Appendix A.

By composing the functors \dagger and \ast , one obtains the functor $_* = \ast \circ \dagger : \mathcal{A} \to \mathcal{A}$ $(a_* = a^*, f_* = (f^{\dagger})^*)$. For a strict dagger compact closed category \mathcal{A} , the functor $_*$ satisfies

$$f_{**} = f$$
, $(f_*)^* = (f^*)_*$.

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A zero-object is an object which is both initial and terminal. For a category with a zero-object 0, there is a composite $0_{a,b}$: $a \to 0 \to b$ for every pair a, b of its objects, and for every other zero-object 0' of this category, the composite $a \to 0' \to b$ is equal to $0_{a,b}$. A *biproduct* of a_1 and a_2 in a category with a zero-object consists of a coproduct and a product diagram

$$a_{1} \xrightarrow{\iota^{1}} a_{1} \oplus a_{2} \xleftarrow{\iota^{2}} a_{2}, \qquad a_{1} \xleftarrow{\pi^{1}} a_{1} \oplus a_{2} \xrightarrow{\pi^{2}} a_{2}$$
for which
$$\pi^{j} \circ \iota^{i} = \begin{cases} \mathbf{1}_{a_{i}}, & i = j, \\ \mathbf{0}_{a_{i},a_{j}}, & \text{otherwise}, \end{cases}$$
(2.1)

where $i, j \in \{1, 2\}$ (cf. the equalities A.16–A.17 in Appendix A). For arrows $f_1: a_1 \to c$ and $f_2: a_2 \to c$, the unique arrow $h: a_1 \oplus a_2 \to c$ for which $h \circ \iota^i = f_i, i \in \{1, 2\}$ is denoted by $[f_1, f_2]$, and for arrows $g_1: c \to a_1$ and $g_2: c \to a_2$, the unique arrow $h: c \to a_1 \oplus a_2$ for which $\pi^i \circ h = g_i, i \in \{1, 2\}$ is denoted by $\langle f_1, f_2 \rangle$.

More generally, a biproduct of a family of objects $\{a_j \mid j \in J\}$ consists of a universal cocone (coproduct diagram) and a universal cone (product diagram)

$$\{\iota^j \colon a_j \to \bigoplus_{j \in J} a_j \mid j \in J\}, \qquad \{\pi^j \colon \bigoplus_{j \in J} a_j \to a_j \mid j \in J\}$$

for which the equality 2.1 holds for all $i, j \in J$. A *category with biproducts* is a category with a zero-object and biproducts for every pair of objects. A biproduct is a *dagger biproduct* when for every pair a, b of objects the equalities A.24 of Appendix A hold.

For $f, g: a \to b$ in a category with biproducts whose *codiagonal* and *diagonal* maps are $\mu_b: b \oplus b \to b$ and $\bar{\mu}_a: a \to a \oplus a$, one defines f + g as $\mu_b \circ (f \oplus g) \circ \bar{\mu}_a$. This operation on the set Hom (a, b) of arrows from a to b is commutative and has $0_{a,b}$ as neutral. Moreover, the composition distributes over +. Hence, every category with biproducts may be conceived as a category enriched over the category **Cmd** of *commutative monoids*.

Alternatively, to define biproducts in a category enriched over **Cmd**, it suffices to assume the existence of a bifunctor \oplus , a special object 0, and for every pair of objects a, b the arrows $\pi_{a,b}^1: a \oplus b \to a, \pi_{a,b}^2: a \oplus b \to b$, $\iota_{a,b}^1: a \to a \oplus b$ and $\iota_{a,b}^2: b \to a \oplus b$, for which the equalities A.14–A.19 of Appendix A hold. As a justification of this approach, see the proof of Corollary 1 below.

In a compact closed category with biproducts, tensor distributes over \oplus , i.e., there exist *distributivity isomorphisms* $\tau_{a,b,c}$: $a \otimes (b \oplus c) \rightarrow (a \otimes b) \oplus (a \otimes c)$ and $\upsilon_{a,b,c}$: $(a \oplus b) \otimes c \rightarrow (a \otimes c) \oplus (b \otimes c)$ explicitly given by

$$\tau_{a,b,c} = \langle \mathbf{1}_a \otimes \pi_{b,c}^1, \mathbf{1}_a \otimes \pi_{b,c}^2 \rangle, \qquad \tau_{a,b,c}^{-1} = [\mathbf{1}_a \otimes \iota_{b,c}^1, \mathbf{1}_a \otimes \iota_{b,c}^2],$$
(2.2)

$$\upsilon_{a,b,c} = \langle \pi_{a,b}^1 \otimes \mathbf{1}_c, \pi_{a,b}^2 \otimes \mathbf{1}_c \rangle, \qquad \upsilon_{a,b,c}^{-1} = [\iota_{a,b}^1 \otimes \mathbf{1}_c, \iota_{a,b}^2 \otimes \mathbf{1}_c].$$
(2.3)

(We are aware that it is hard to distinguish between the Latin letter v, which is reserved for the isomorphism from I^* to I and the Greek letter v denoting the isomorphism of the form $(a \oplus b) \otimes c \rightarrow (a \otimes c) \oplus (b \otimes c)$, but we decided to follow the notation from Ref. [20] relevant for the strict compact closed structure, and from Ref. [1] which is relevant for categorical quantum protocols.)

In a compact closed categories with biproducts, the *scalars*, i.e., the endomorphisms from I to I form a commutative semiring Hom (I, I). The multiplication in this semiring is given by composition, for which $\mathbf{1}_I$ is the neutral, and the addition is defined as above. (We will omit \circ when we compose, i.e., multiply, scalars.) For a scalar $s: I \to I$ and an object a of such a category, one defines the arrow $s_a: a \to a$ as the composition

$$a \xrightarrow{\lambda_a^{-1}} I \otimes a \xrightarrow{s \otimes a} I \otimes a \xrightarrow{\lambda_a} a,$$

and the operation $s \bullet$ on arrows such that for $f : a \to b$, the arrow $s \bullet f$ is $f \circ s_a$. It is straightforward to check that this new operation satisfies the following equalities.

$$a \otimes (s \bullet f) = s \bullet (a \otimes f), \quad (s \bullet f) \otimes a = s \bullet (f \otimes a), \tag{2.4}$$

$$(s_2 \bullet f_2) \circ (s_1 \bullet f_1) = s_2 s_1 \bullet (f_2 \circ f_1), \tag{2.5}$$

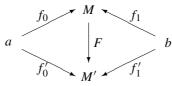
$$\langle s \bullet f_1, \dots, s \bullet f_n \rangle = s \bullet \langle f_1, \dots, f_n \rangle.$$
 (2.6)

Example 1 As a paradigm for dagger compact closed category with dagger biproducts, we use the category **fdHilb** of finite dimensional Hilbert spaces over the field \mathbb{C} of complex numbers. The objects of this category are finite dimensional Hilbert spaces (finite dimensional vector spaces with inner product). The arrows of this category correspond to (bounded) linear maps between vector spaces. Dagger is given by the adjoint map. Since every vector space over \mathbb{C} of dimension *n* is isomorphic to \mathbb{C}^n , we can pass from **fdHilb** to its skeleton consisting of objects of the form \mathbb{C}^n . By choosing orthogonal bases of such objects, the linear maps are envisaged as matrices. In this case, dagger corresponds to the usual adjoint of matrices (conjugation and transposition), and the operation * on arrows corresponds to the complex conjugation of an operator (matrix). Also, the operation * on arrows is given by a matrix transpose.

3 The category 1Cob

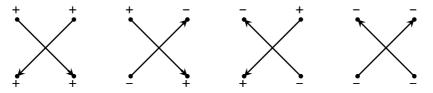
The category 1**Cob** of one-dimensional cobordisms has as objects closed oriented zero-dimensional manifolds, i.e., finite (possibly empty) sequences of points together with their orientation (either + or –). For example, an object of 1**Cob** is ++-+--. Since there will be several roles of \emptyset in this paper, we denote the empty sequence of points by o.

A compact oriented one-dimensional topological manifold with boundary, i.e., a finite collection of oriented circles and line segments, is called here 1-manifold. For objects a and b of 1**Cob**, a 1-cobordism from a to b is a triple $(M, f_0: a \to M, f_1: b \to M)$, where M is a 1-manifold and f_0, f_1 are embeddings. The boundary of M is $\Sigma_0 \coprod \Sigma_1$ and its orientation is induced from the orientation of M (the initial point of an oriented segment is + while the terminal is –). The embedding f_0 is orientation preserving and its image is Σ_0 , while the embedding f_1 is orientation reversing and its image is Σ_1 . Two cobordisms (M, f_0, f_1) and (M', f'_0, f'_1) from a to b are equivalent, when there is an orientation preserving homeomorphism $F: M \to M'$ such that the following diagram commutes.



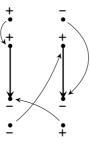
The equivalence classes of 1-cobordisms are the *arrows* of 1**Cob**. The identity $\mathbf{1}_a: a \to a$ is the equivalence class of $(a \times I, x \mapsto (x, 0), x \mapsto (x, 1))$, which in the case a = o stands for the empty 1-cobordism from the empty sequence of oriented points o to itself. Two cobordisms $(M, f_0, f_1): a \to b$ and $(N, g_0, g_1): b \to c$ are composed by "gluing", i.e., by making the pushout of $M \xleftarrow{f_1} b \xrightarrow{g_0} N$. All the arrows of 1**Cob** are illustrated so that the source of an arrow is at the top, while its target is at the bottom of the picture. Therefore, the direction of pictures is *top to bottom*, a convention used in Ref. [29]. Note that some authors use a different convection, *left to right*, or *bottom to top* [22,34]. The latter is presumably the most popular in the physics literature.

The category 1**Cob** is dagger strict compact closed. We have symmetric monoidal structure on 1**Cob** in which \otimes on objects is defined by concatenation, the empty sequence *o* is the neutral and serves as the unit object *I*, while \otimes on arrows is given by putting two cobordisms "side by side". The arrows α and λ are identities and symmetry σ is generated by transpositions $++ \rightarrow ++, +- \rightarrow -+, -+ \rightarrow +-$ and $-- \rightarrow --$. These transpositions are illustrated as follows:

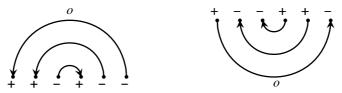


For example, the transposition $+- \rightarrow -+$ is a cobordism given by the manifold consisting of two oriented segments and two embeddings of the source +- and the target -+ into its boundary (when a sign is mapped to the same

sign, then the boundary point belongs to the source, and it is mapped to the opposite sign, then the boundary point belongs to the target of the cobordism).

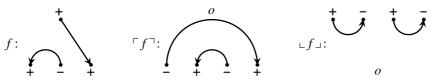


The dual a^* of an object a is the reversed sequence of points with reversed orientation. For example, if a = +--, then $a^* = ++-$. (Note that this definition differs from the one given in Ref. [29] where just the orientation was reversed—both definitions are correct in presence of symmetry.) The arrows $\eta: o \to a^* \otimes a$ and $\varepsilon: a \otimes a^* \to o$, for a as above are the following cobordisms:



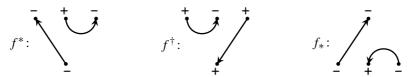
It is not difficult to check that the arrows $u_{a,b}$, v and w_a are all identities.

Let $f: a \to b$ be an arrow of 1**Cob** represented by a triple $(M, f_0: a \to M, f_1: b \to M)$. Its name $\lceil f \rceil: o \to a^* \otimes b$ is represented by the triple $(M, g_0: o \to M, g_1: a^* \otimes b \to M)$, where for every point x of a and the corresponding point \bar{x} of a^* , we have $g_1(\bar{x}) = f_0(x)$ and for every point y of b, we have $g_1(y) = f_1(y)$. The coname $\lfloor f \rfloor$ of f is defined in 1**Cob** analogously.



The arrow $f^*: b^* \to a^*$ is represented by the triple $(M, h_0: b^* \to M, h_1: a^* \to M)$, where for every point x of a and the corresponding point \bar{x} of a^* , we have $h_1(\bar{x}) = f_0(x)$, and for every point y of b and the corresponding point \bar{y} of b^* , we have $h_0(\bar{y}) = f_1(y)$.

The cobordism $f^{\dagger}: b \to a$ is obtained by reversing the orientation of the 1-manifold representing the cobordism $f: a \to b$. It is not hard to check that the equalities A.21–A.23 of Appendix A hold.



By the above definitions of f^* and f^{\dagger} for a cobordism $f: a \to b$, it is straightforward to reconstruct the cobordism $f_* = (f^{\dagger})^*: a^* \to b^*$.

4 A pair of free categories

We start with a construction of a dagger compact closed category \mathcal{F}^{\dagger} with dagger biproducts freely generated by a single object p and a set Γ of unitary endomorphisms on this object. An arrow $f: a \to b$ in a dagger category

is *unitary* when $f^{\dagger}: b \to a$ is its both-sided inverse. The universal property of \mathcal{F}^{\dagger} is the following: for every dagger compact closed category \mathcal{C} with dagger biproducts, and a function φ from the set Γ to the set of unitary endomorphisms of an object c of \mathcal{C} , there exists a unique functor $F: \mathcal{F}^{\dagger} \to \mathcal{C}$ strictly preserving the whole structure, such that Fp = c and for every $\gamma \in \Gamma$, $F\gamma = \varphi(\gamma)$.

Our construction of this category is syntactical; it is akin to the construction of the category FA from Ref. [20, Sections 3–4], and it follows the construction of the category \mathcal{F}_P from Ref. [29, Section 4]. As noted in [20], it is "perfectly general, applying to categories with any explicitly given equational extra structure". The *objects* of \mathcal{F}^{\dagger} are the formulae built out of a single letter p and the constants I and 0, with the help of one unary connective * (written as a superscript) and two binary connectives \otimes and \oplus .

The arrows of \mathcal{F}^{\dagger} are obtained as equivalence classes of *terms* built in the following manner. We start with *primitive terms*, which are of the form γ , γ^{-1} for every $\gamma \in \Gamma$, or $\mathbf{1}_a$, $\alpha_{a,b,c}$, $\alpha_{a,b,c}^{-1}$, λ_a , λ_a^{-1} , $\sigma_{a,b}$, η_a , ε_a , $\pi_{a,b}^1$, $\pi_{a,b}^2$, $\iota_{a,b}^1$, $\iota_{a,b}^2$

$$\gamma \circ \gamma^{-1} = \mathbf{1}_p = \gamma^{-1} \circ \gamma,$$

$$\gamma^{\dagger} = \gamma^{-1}.$$

$$(4.1)$$

$$(4.2)$$

On the other hand, consider the category \mathcal{F} with the same objects and the same primitive terms as \mathcal{F}^{\dagger} , just the terms of \mathcal{F} are constructed without the unary operational symbol \dagger . The *arrows* of \mathcal{F} , are the equivalence classes of these terms, modulo the congruence generated by the equalities A.1–A.19 and 4.1. The category \mathcal{F} is a compact closed category with biproducts freely generated by the group (envisaged as a category with one object) freely generated by the set Γ . The universal property of \mathcal{F} is the following: for every compact closed category \mathcal{C} with biproducts, and a function φ from the set Γ to the set of automorphisms of an object c of \mathcal{C} , there exists a unique functor $F: \mathcal{F} \to \mathcal{C}$ strictly preserving the whole structure, such that Fp = c and for every $\gamma \in \Gamma$, $F\gamma = \varphi(\gamma)$.

Proposition 1 *The categories* \mathcal{F}^{\dagger} *and* \mathcal{F} *are isomorphic.*

Proof From the equalities A.20–A.26, it follows that every arrow of \mathcal{F}^{\dagger} (as an equivalence class of terms) contains a \dagger -free term. Also, every equality assumed for \mathcal{F}^{\dagger} in which \dagger appears boils down to the trivial identity after \dagger elimination at both sides. Thus, the identity on objects and the function on arrows that maps the equivalence class
of a term in \mathcal{F} to the equivalence class of the same term in \mathcal{F}^{\dagger} is an isomorphism between these two categories. \Box

5 Injections and projections

For the functor *: $\mathcal{F}^{op} \to \mathcal{F}$ defined as in Sect. 2, the unit η and the counit ε become *dinatural*, i.e., for $f: a \to b$ the following equalities hold:

$$(a^* \otimes f) \circ \eta_a = (f^* \otimes b) \circ \eta_b, \qquad \varepsilon_a \circ (a \otimes f^*) = \varepsilon_b \circ (f \otimes b^*). \tag{5.1}$$

Also, for arrows $f, g : a \to b$ in \mathcal{F} , the following equality holds,

$$(f+g)^* = f^* + g^*.$$
(5.2)

Definition 1 Let *a* be an object of \mathcal{F} . By induction on complexity of *a*, we define two finite sequences $I_a = (\iota_a^0, \ldots, \iota_a^{n-1})$ (the *injections* of *a*) and $\Pi_a = (\pi_a^0, \ldots, \pi_a^{n-1})$ (the *projections* of *a*) of arrows of \mathcal{F} in the following way. If *a* is the letter *p* or either *I* or 0, then n = 1 and $I_a = (\mathbf{1}_a) = \Pi_a$. Let us assume that $I_{a_1} = (\iota_1^0, \ldots, \iota_1^{n_1-1})$, $\Pi_{a_1} = (\pi_1^0, \ldots, \pi_1^{n_1-1})$ and $I_{a_2} = (\iota_2^0, \ldots, \iota_2^{n_2-1})$, $\Pi_{a_2} = (\pi_2^0, \ldots, \pi_2^{n_2-1})$ are already defined. For $\lfloor x \rfloor$ being the *floor* function of a real *x*, i.e., the greatest integer less than or equal to *x*, and *i* mod *n* being the residue of *i* modulo *n*, we have the following.

$$\begin{split} & \text{If } a = a_1 \otimes a_2, \text{ then } n = n_1 \cdot n_2, \text{ and for } 0 \leq i < n_1 \cdot n_2, \\ & \iota_a^i = \iota_1^{\lfloor i/n_2 \rfloor} \otimes \iota_2^{i \mod n_2}, \qquad \pi_a^i = \pi_1^{\lfloor i/n_2 \rfloor} \otimes \pi_2^{i \mod n_2}. \\ & \text{* If } a = a_1^*, \text{ then } n = n_1, \text{ and for } 0 \leq i < n_1, \\ & \iota_a^i = (\pi_1^i)^*, \quad \pi_a^i = (\iota_1^i)^*. \\ & \oplus \text{ If } a = a_1 \oplus a_2, \text{ then } n = n_1 + n_2, \text{ and for } 0 \leq i < n_1 + n_2, s_i = \left\lfloor \frac{\min\{i, n_1\}}{n_1} \right\rfloor \\ & \iota_a^i = \left\{ \begin{array}{c} \iota_{a_1, a_2}^1 \circ \iota_1^i, & 0 \leq i < n_1, \\ \iota_{a_1, a_2}^2 \circ \iota_2^{i-n_1}, & 0 \leq i < n_1, \\ \iota_{a_1, a_2}^2 \circ \iota_2^{i-n_1}, & 0 \in i < n_1, \\ \eta_a^i = \left\{ \begin{array}{c} \pi_1^i \circ \pi_{a_1, a_2}^1, & 0 \leq i < n_1, \\ \pi_2^{i-n_1} \circ \pi_{a_1, a_2}^2, & 0 \leq i < n_1, \\ \pi_2^{i-n_1} \circ \pi_{a_1, a_2}^2, & 0 \leq i < n_1, \\ \pi_2^{i-n_1} \circ \pi_{a_1, a_2}^2, & 0 \in i < n_1, \\ \pi_2^{i-n_1} \circ \pi_{a_1, a_2}^2, & 0 \in i < n_1, \\ \pi_2^{i-n_1} \circ \pi_{a_1, a_2}^2, & 0 \in i < n_1, \\ \pi_2^{i-n_1} \circ \pi_{a_1, a_2}^2, & 0 \in i < n_1, \\ \pi_2^{i-n_1} \circ \pi_{a_1, a_2}^2, & 0 \in i < n_1, \\ \pi_2^{i-n_1} \circ \pi_{a_1, a_2}^2, & 0 \in i < n_1, \\ \pi_2^{i-n_1} \circ \pi_{a_1, a_2}^2, & 0 \in i < n_1, \\ \pi_2^{i-n_1} \circ \pi_{a_1, a_2}^2, & 0 \in i < n_1, \\ \pi_2^{i-n_1} \circ \pi_{a_1, a_2}^2, & 0 \in i < n_1, \\ \pi_2^{i-n_1} \circ \pi_{a_1, a_2}^2, & 0 \in i < n_1, \\ \pi_2^{i-n_1} \circ \pi_{a_1, a_2}^2, & 0 \in i < n_1, \\ \pi_2^{i-n_1} \circ \pi_{a_1, a_2}^2, & 0 \in i < n_1, \\ \pi_2^{i-n_1} \circ \pi_{a_1, a_2}^2, & 0 \in i < n_1, \\ \pi_2^{i-n_1} \circ \pi_{a_1, a_2}^2, & 0 \in i < n_1, \\ \pi_2^{i-n_1} \circ \pi_{a_1, a_2}^2, & 0 \in i < n_1, \\ \pi_2^{i-n_1} \circ \pi_{a_1, a_2}^2, & 0 \in i < n_1, \\ \pi_2^{i-n_1} \circ \pi_{a_1, a_2}^2, & 0 \in i < n_1, \\ \pi_2^{i-n_1} \circ \pi_{a_1, a_2}^2, & 0 \in i < n_1, \\ \pi_2^{i-n_1} \circ \pi_{a_1, a_2}^2, & 0 \in i < n_1, \\ \pi_2^{i-n_1} \circ \pi_{a_1, a_2}^2, & 0 \in i < n_1, \\ \pi_2^{i-n_1} \circ \pi_2^{i-n_1} \circ \pi_2^2, & 0 \in i < n_1, \\ \pi_2^{i-n_1} \circ \pi_2^{i-n_2} \circ \pi_2^2, & 0 \in i < n_1 \\ \pi_2^{i-n_1} \circ \pi_2^{i-n_1} \circ \pi_2^2, & 0 \in i < n_1 \\ \pi_2^{i-n_1} \circ \pi_2^{i-n_1} \circ \pi_2^2, & 0 \in i < n_1 \\ \pi_2^{i-n_1} \circ \pi_2^{i-n_1} \circ \pi_2^2, & 0 \in i < n_1 \\ \pi_2^{i-n_1} \circ \pi$$

Example 2 Let $a = (p \oplus I) \oplus 0$ and $b = ((p \oplus 0) \oplus p) \otimes (I \oplus p)^*$. Then ι_a^i and π_a^i for $0 \le i < 3$ as well as ι_b^j and π_b^j for $0 \le j < 6$ are given in the following tables.

ι_a^0	$\iota^1_{p\oplus I,0}\circ\iota^1_{p,I}$	l.	ι_b^0	$(\iota^1_{p\oplus 0,p}\circ\iota^1_{p,0})\otimes (\pi^1_{I,p})^*$	π_b^0	$(\pi^1_{p,0}\circ\pi^1_{p\oplus 0,p})\otimes(\iota^1_{I,p})^*$
ι^1_a	$\iota^1_{p\oplus I,0}\circ\iota^2_{p,I}$	L	ι_b^1	$(\iota^1_{p\oplus 0,p}\circ\iota^1_{p,0})\otimes (\pi^2_{I,p})^*$	π_b^1	$(\pi^1_{p,0}\circ\pi^1_{p\oplus 0,p})\otimes(\iota^2_{I,p})^*$
ι_a^2	$\iota^2_{p\oplus I,0}$	l	ι_b^2	$(\iota^1_{p\oplus 0,p}\circ\iota^2_{p,0})\otimes (\pi^1_{I,p})^*$	π_b^2	$(\pi_{p,0}^2\circ\pi_{p\oplus 0,p}^1)\otimes(\iota_{I,p}^1)^*$
π_a^0	$\pi^1_{p,I} \circ \pi^1_{p \oplus I,0}$	L	ι_b^3	$(\iota^1_{p\oplus 0,p}\circ\iota^2_{p,0})\otimes (\pi^2_{I,p})^*$	π_b^3	$(\pi_{p,0}^2\circ\pi_{p\oplus 0,p}^1)\otimes(\iota_{I,p}^2)^*$
π^1_a	$\pi_{p,I}^2\circ\pi_{p\oplus I,0}^1$	Ľ	ι_b^4	$\iota^2_{p\oplus 0,p}\otimes (\pi^1_{I,p})^*$	π_b^4	$\pi^2_{p\oplus 0,p}\otimes (\iota^1_{I,p})^*$
π_a^2	$\pi^2_{p\oplus I,0}$	L.	ι_b^5	$\iota^2_{p\oplus 0,p}\otimes (\pi^2_{I,p})^*$	π_b^5	$\pi^2_{p\oplus 0,p}\otimes (\iota^2_{I,p})^*$

Remark 1 For every $0 \le i < n$, the target of ι_a^i and the source of π_a^i are both equal to a, while the source a^i of ι_a^i is equal to the target of π_a^i , and a^i is \oplus -free. Moreover, if a is \oplus -free, then $I_a = (\mathbf{1}_a) = \Pi_a$.

The following proposition establishes the desired properties of injections and projections.

Proposition 2 For every object a of \mathcal{F}

$$\pi_a^j \circ \iota_a^i = \begin{cases} \mathbf{I}_{a^i}, & i = j, \\ \mathbf{0}_{a^i, a^j}, & \text{otherwise,} \end{cases} \qquad \sum_{i=0}^{n-1} \iota_a^i \circ \pi_a^i = \mathbf{I}_a.$$

Proof We proceed by induction on complexity of a. When a is p, I or 0, all injections and projections are identities, and the claim holds. For the inductive step, we consider the following three cases.

(1) Suppose that $a = a_1 \otimes a_2$, where $|I_{a_1}| = |\Pi_{a_1}| = n_1$ and $|I_{a_2}| = |\Pi_{a_2}| = n_2$. Then we have

$$\begin{aligned} \pi_a^j \circ \iota_a^i &= \left(\pi_1^{\lfloor \frac{j}{n_2} \rfloor} \otimes \pi_2^{j \mod n_2}\right) \circ \left(\iota_1^{\lfloor \frac{i}{n_2} \rfloor} \otimes \iota_2^{i \mod n_2}\right) \\ &= \left(\pi_1^{\lfloor \frac{j}{n_2} \rfloor} \circ \iota_1^{\lfloor \frac{i}{n_2} \rfloor}\right) \otimes \left(\pi_2^{j \mod n_2} \circ \iota_2^{i \mod n_2}\right), \end{aligned}$$

and the first claim follows by the inductive hypothesis. Also, using the inductive hypothesis and the equality A.27 of Appendix A, we have

$$\sum_{i=0}^{n-1} \iota_a^i \circ \pi_a^i = \sum_{i=0}^{n-1} \left(\iota_1^{\lfloor \frac{i}{n_2} \rfloor} \otimes \iota_2^{i \mod n_2} \right) \circ \left(\pi_1^{\lfloor \frac{i}{n_2} \rfloor} \otimes \pi_2^{i \mod n_2} \right)$$

$$=\sum_{i=0}^{n-1} \left(\iota_{1}^{\lfloor \frac{i}{n_{2}} \rfloor} \circ \pi_{1}^{\lfloor \frac{i}{n_{2}} \rfloor} \right) \otimes \left(\iota_{2}^{i \mod n_{2}} \circ \pi_{2}^{i \mod n_{2}} \right)$$

$$=\sum_{k=0}^{n_{1}-1} \left((\iota_{1}^{k} \circ \pi_{1}^{k}) \otimes \sum_{l=0}^{n_{2}-1} (\iota_{2}^{l} \circ \pi_{2}^{l}) \right) = \sum_{k=0}^{n_{1}-1} (\iota_{1}^{k} \circ \pi_{1}^{k}) \otimes a_{2}$$

$$= \left(\sum_{k=0}^{n_{1}-1} (\iota_{1}^{k} \circ \pi_{1}^{k}) \right) \otimes a_{2} = \mathbf{1}_{a_{1}} \otimes a_{2} = \mathbf{1}_{a_{1} \otimes a_{2}}.$$

(2) When $a = a_1^*$, we have

$$\pi_a^j \circ \iota_a^i = (\iota_1^j)^* \circ (\pi_1^i)^* = (\pi_1^i \circ \iota_1^j)^*,$$

and the first claim follows by the inductive hypothesis. Also, using the inductive hypothesis and the equality 5.2, we have

$$\sum_{i=0}^{n-1} \iota_a^i \circ \pi_a^i = \sum_{i=0}^{n-1} (\pi_1^i)^* \circ (\iota_1^i)^* = \sum_{i=0}^{n-1} (\iota_1^i \circ \pi_1^i)^* = \left(\sum_{i=0}^{n-1} \iota_1^i \circ \pi_1^i\right)^* = (\mathbf{1}_{a_1})^* = \mathbf{1}_a$$

(3) Suppose that $a = a_1 \oplus a_2$, and again $|I_{a_1}| = |\Pi_{a_1}| = n_1$ and $|I_{a_2}| = |\Pi_{a_2}| = n_2$. We have $\pi_a^j \circ \iota_a^i = \pi_{1+s_j}^{j-n_1 \cdot s_j} \circ \pi_{a_1,a_2}^{1+s_j} \circ \iota_{1+s_i}^{i-n_1 \cdot s_i}$. Since $i \neq j$ implies $1 + s_i \neq 1 + s_j$ or $i - n_1 \cdot s_i \neq j - n_1 \cdot s_j$, the first claim follows according to the inductive hypothesis. For the second claim, we have

$$\begin{split} \sum_{i=0}^{n-1} \iota_a^i \circ \pi_a^i &= \sum_{i=0}^{n-1} \iota_{a_1,a_2}^{1+s_i} \circ \iota_{1+s_i}^{i-n_1 \cdot s_i} \circ \pi_{1+s_i}^{i-n_1 \cdot s_i} \circ \pi_{a_1,a_2}^{1+s_i} \\ &= \sum_{i=0}^{n_1-1} \iota_{a_1,a_2}^1 \circ \iota_1^i \circ \pi_1^i \circ \pi_{a_1,a_2}^1 + \sum_{j=0}^{n_2-1} \iota_{a_1,a_2}^2 \circ \iota_2^j \circ \pi_2^j \circ \pi_{a_1,a_2}^2 \\ &= \iota_{a_1,a_2}^1 \circ \left(\sum_{i=0}^{n_1-1} \iota_1^i \circ \pi_1^i \right) \circ \pi_{a_1,a_2}^1 + \iota_{a_1,a_2}^2 \circ \left(\sum_{j=0}^{n_2-1} \iota_2^j \circ \pi_2^j \right) \circ \pi_{a_1,a_2}^2 \\ &= \iota_{a_1,a_2}^1 \circ \pi_{a_1,a_2}^1 + \iota_{a_1,a_2}^2 \circ \pi_{a_1,a_2}^2 = \mathbf{1}_{a_1 \oplus a_2}. \end{split}$$

As a corollary of Proposition 2, we have the following.

Corollary 1 For every object a of \mathcal{F} , the cocone (a, I_a) together with the cone (a, Π_a) make a biproduct.

Proof To show that the cocone (a, I_a) is universal, consider for $0 \le i < n$ arrows $f^i : a^i \to c$ of \mathcal{F} and define $h: a \to c$ to be $\sum_{i=0}^{n-1} f^i \circ \pi_a^i$. For every $0 \le i < n$, by the left-hand side equality of Proposition 2, we have that h satisfies $h \circ \iota_a^i = f^i$. Assume that $h': a \to c$ for every $0 \le i < n$ satisfies $h' \circ \iota_a^i = f^i$. We conclude that

$$h' \circ \sum_{i=0}^{n-1} \iota_a^i \circ \pi_a^i = \sum_{i=0}^{n-1} f^i \circ \pi_a^i = h,$$

and by the right-hand side equality of Proposition 2, we have h' = h. That (a, Π_a) is a universal cone is proved analogously.

6 A normalization

Our normalization of arrows of the category \mathcal{F} is a procedure derived from the one developed in Ref. [29, Section 5]. The goal is to represent every arrow of \mathcal{F} , whose source and target are \oplus -free, by a term free of occurrences of \oplus , ι and π .

For every arrow $u: a \to b$ of \mathcal{F} , where $I_a = (\iota_a^0, \ldots, \iota_a^{n-1})$, $\Pi_b = (\pi_b^0, \ldots, \pi_b^{m-1})$, let M_u be the $m \times n$ matrix whose ij-entry is $\pi_b^i \circ u \circ \iota_a^j: a^j \to b^i$. Let X be an $m \times n$ matrix whose ij-entry is an arrow of \mathcal{F} from a^j to b^i and let Y be a $q \times r$ matrix whose ij-entry is an arrow of \mathcal{F} from c^j to d^i . We define $X \otimes Y$ as the Kronecker product of matrices over a field, save that the multiplication in the field is replaced by the tensor product of arrows in \mathcal{F} . For example,

$$\begin{pmatrix} x_{00} & x_{01} & x_{02} \\ x_{10} & x_{11} & x_{12} \end{pmatrix} \otimes \begin{pmatrix} y_{00} & y_{01} \\ y_{10} & y_{11} \end{pmatrix}$$

is

 $\begin{pmatrix} x_{00} \otimes y_{00} & x_{00} \otimes y_{01} & x_{01} \otimes y_{00} & x_{01} \otimes y_{01} & x_{02} \otimes y_{00} & x_{02} \otimes y_{01} \\ x_{00} \otimes y_{10} & x_{00} \otimes y_{11} & x_{01} \otimes y_{10} & x_{01} \otimes y_{11} & x_{02} \otimes y_{10} & x_{02} \otimes y_{11} \\ x_{10} \otimes y_{00} & x_{10} \otimes y_{01} & x_{11} \otimes y_{00} & x_{11} \otimes y_{01} & x_{12} \otimes y_{00} & x_{12} \otimes y_{01} \\ x_{10} \otimes y_{10} & x_{10} \otimes y_{11} & x_{11} \otimes y_{10} & x_{11} \otimes y_{11} & x_{12} \otimes y_{10} & x_{12} \otimes y_{11} \end{pmatrix}$

Also, we define $X \oplus Y$ as the matrix of arrows of \mathcal{F} , schematically presented as

 $\left(\begin{array}{cc} X & 0 \\ 0 & Y \end{array}\right).$

More precisely, for $X = (x_{ij})_{m \times n}$ and $Y = (y_{ij})_{q \times r}$ as above, $X \oplus Y$ is the $(m + q) \times (n + r)$ matrix whose *ij*-entry is:

- 1. x_{ij} , when i < m, j < n,
- 2. $y_{(i-m)(j-n)}$, when $i \ge m, j \ge n$,
- 3. $0_{a^{j}, d^{i-m}}$, when $i \ge m, j < n$,
- 4. $0_{c^{j-n},b^i}$, when $i < m, j \ge n$.

If m = q, n = r, for every $0 \le j < n$, $a^j = c^j$, and for every $0 \le i < m$, $b^i = d^i$, i.e., X and Y are of the same type having the corresponding elements in the same hom-sets, then X + Y is the matrix of the same type whose ij-entry is $x_{ij} + y_{ij}$. If n = q and for every $0 \le k < n$, $a^k = d^k$, i.e., X is an $m \times n$ matrix, Y is an $n \times r$ matrix and for every $0 \le i < m$, $0 \le j < r$, $0 \le k < n$ the composition $x_{ik} \circ y_{kj}$ is defined, then we define $X \circ Y$ as the $m \times r$ matrix whose ij-entry is $\sum_{k=0}^{n-1} x_{ik} \circ y_{kj}$ (this sum is defined since every $x_{ik} \circ y_{kj}$ is from c^j to b^i).

Just by omitting the case (2) of Ref. [29, Proposition 5.1] we obtain the following.

Proposition 3 *For* $\bullet \in \{\otimes, \oplus, +, \circ\}$ *, we have*

$$M_{u_1 \bullet u_2} = M_{u_1} \bullet M_{u_2}$$

Our next proposition is related to Ref. [29, Propositions 5.2, 8.2]

Proposition 4 If u is a primitive term of \mathcal{F} , then all the entries of the matrix M_u are primitive terms of \mathcal{F} , not of the form π or ι , whose indices are \oplus -free.

Proof We illustrate just a couple of cases. If u is γ , for $\gamma \in \Gamma$, then M_u is a 1×1 matrix whose only entry is γ . The same holds when γ is replaced by γ^{-1} . If u is $\alpha_{a,b,c}^{-1}$, then for some i_1, i_2, i_3 and j_1, j_2, j_3

$$(M_u)_{i,j} = \pi^i_{a\otimes(b\otimes c)} \circ \alpha^{-1}_{a,b,c} \circ \iota^j_{(a\otimes b)\otimes c}$$

= $(\pi^{i_1}_a \otimes (\pi^{i_2}_b \otimes \pi^{i_3}_c)) \circ \alpha^{-1}_{a,b,c} \circ ((\iota^{j_1}_a \otimes \iota^{j_2}_b) \otimes \iota^{j_3}_c)$

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$$=\begin{cases} \alpha_{a^{i_1},b^{i_2},c^{i_3}}^{-1}, & i_1 = j_1, \ i_2 = j_2, \ i_3 = j_3, \\ 0_{(a^{j_1} \otimes b^{j_2}) \otimes c^{j_3},a^{i_1} \otimes (b^{i_2} \otimes c^{i_3}), \\ \end{cases} \text{ otherwise.}$$

If *u* is $\sigma_{a,b}$, then for some i_1, i_2 and j_1, j_2

$$(M_u)_{i,j} = \pi_{b\otimes a}^i \circ \sigma_{a,b} \circ \iota_{a\otimes b}^j$$

= $(\pi_b^{i_1} \otimes \pi_a^{i_2}) \circ \sigma_{a,b} \circ (\iota_a^{j_1} \otimes \iota_b^{j_2})$
=
$$\begin{cases} \sigma_{a^{j_1},b^{j_2}}, & j_1 = i_2, \ j_2 = i_1, \\ 0_{a^{j_1} \otimes b^{j_2},b^{j_1} \otimes a^{j_2}}, & \text{otherwise.} \end{cases}$$

If u is ε_a , then M_u is a row matrix and for some j_1 , j_2 we have

$$(M_u)_{1,j} = \varepsilon_a \circ \iota_{a\otimes a^*}^j = \varepsilon_a \circ (\iota_a^{j_1} \otimes (\pi_a^{j_2})^*) \stackrel{(5.1)}{=} \varepsilon_{a^{j_2}} \circ ((\pi_a^{j_2} \circ \iota_a^{j_1}) \otimes (a^{j_2})^*)$$
$$= \begin{cases} \varepsilon_{a^{j_2}}, & j_1 = j_2, \\ 0_{a^{j_1} \otimes (a^{j_2})^*, I}, & \text{otherwise.} \end{cases}$$

If u is $\pi_{a,b}^1$, then $(M_u)_{i,j} = \pi_a^i \circ \pi_{a,b}^1 \circ \iota_{a\oplus b}^j$, which is either $\pi_a^i \circ \pi_{a,b}^1 \circ \iota_{a,b}^1 \circ \iota_{a}^{j_1}$ for some j_1 , or $\pi_a^i \circ \pi_{a,b}^1 \circ \iota_{a,b}^{j_2} \circ \iota_{b}^{j_2}$, for some j_2 . The statement holds since

$$\pi_{a}^{i} \circ \pi_{a,b}^{1} \circ \iota_{a,b}^{1} \circ \iota_{a}^{j_{1}} = \begin{cases} \mathbf{1}_{a^{i}}, & i = j_{1}, \\ \mathbf{0}_{a^{j_{1}},a^{i}}, & \text{otherwise,} \end{cases} \qquad \pi_{a}^{i} \circ \pi_{a,b}^{1} \circ \iota_{a,b}^{2} \circ \iota_{b}^{j_{2}} = \mathbf{0}_{b^{j_{2}},a^{i}}.$$

We proceed analogously when *u* is $\mathbf{1}_a$, $\alpha_{a,b,c}$, λ_a , λ_a^{-1} , η_a , $\pi_{a,b}^2$, $\iota_{a,b}^1$, $\iota_{a,b}^2$ or $\mathbf{0}_{a,b}$.

Corollary 2 For every arrow u of \mathcal{F} , every entry of M_u is expressible free of \oplus , ι and π .

As a consequence of Remark 1 and Corollary 2, we have the following.

Corollary 3 Every arrow of \mathcal{F} whose source and target are \oplus -free is expressible free of \oplus , ι , and π .

7 The category $1Cob_{\mathfrak{G}}^{\oplus}$ and coherence

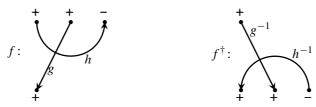
The aim of this section is to introduce a category providing a diagrammatical checking of the validity of quantum protocols. We start with a set Γ (usually finite and non-empty) and a group \mathfrak{G} freely generated by Γ . The category 1**Cob** and the group \mathfrak{G} deliver the category 1**Cob** through the following construction. The objects of 1**Cob** are the objects of 1**Cob** and to define the arrows of 1**Cob**, we introduce the notions of \mathfrak{G} -components and \mathfrak{G} -cobordisms first.

A \mathfrak{G} -component is a connected, oriented 1-manifold, possibly with boundaries, together with an element of \mathfrak{G} . When a \mathfrak{G} -component is closed, we call it \mathfrak{G} -circle, otherwise it is a \mathfrak{G} -segment. We call the element of \mathfrak{G} associated to a component the *label* of this component.

A \mathfrak{G} -cobordism from a to b is a finite collection of \mathfrak{G} -components whose underlying manifold is M, together with two embeddings $f_0: a \to M$ and $f_1: b \to M$ such that (M, f_0, f_1) is a 1-cobordism from a to b. Two \mathfrak{G} -cobordisms are equivalent, when the underlying 1-cobordisms are equivalent and the homeomorphism F witnessing this equivalence satisfies:

- 1. Every segment and its F-image are labeled by the same element of \mathfrak{G} ;
- 2. The labels of a circle and its *F*-image could differ only in a circular permutation, i.e., if one is of the form $g_2 \cdot g_1$, the other could be $g_1 \cdot g_2$.

The operation \dagger on \mathfrak{G} -cobordisms is defined so that it is applied to the underlying cobordism and every label is replaced by its inverse.



The category 1**Cob**_{\mathfrak{G}} has the equivalence classes of \mathfrak{G} -cobordisms as arrows. The identity $\mathbf{1}_a: a \to a$ is the ordinary identity cobordism in which every segment is labeled by the *neutral* e of \mathfrak{G} . Two \mathfrak{G} -cobordisms are composed so that the underlying 1-cobordisms are composed in the ordinary manner. It remains to label the resulting segments and circles: if the segments l_1, \ldots, l_k with labels g_1, \ldots, g_k respectively, are glued together in a segment or a circle of the resulting 1-cobordism so that the terminal point of l_i is identified with the initial point of l_{i+1} , then $g_k \cdot \ldots \cdot g_1$ is the ("a" in the case of a circle) label of the resulting component. The category 1**Cob**_{\mathfrak{G}} has dagger strict compact closed structure inherited from 1**Cob** (all segments in canonical arrows σ , η and ε are labeled in 1**Cob**_{\mathfrak{G}} by the neutral e of \mathfrak{G}).

Let us compare the above construction with the construction of the category $G\mathcal{A}$ given in Ref. [20], for \mathcal{A} being the groupoid \mathfrak{G} , i.e., the category with a single object p whose arrows are the elements of \mathfrak{G} and the composition is the multiplication in \mathfrak{G} . The main theorem of Ref. [20] claims that $G\mathfrak{G}$ is a compact closed category freely generated by the category \mathfrak{G} . This means that for every compact closed category \mathcal{C} and a function φ from the set Γ to the set of automorphisms of an object c of \mathcal{C} , there exists a unique functor

$F\colon G\mathfrak{G}\to \mathcal{C}$

that strictly preserves compact closed structure, and such that Fp = c and for every $\gamma \in \Gamma$, $F\gamma = \varphi(\gamma)$.

One could easily conclude that $1\operatorname{Cob}_{\mathfrak{G}}$ is a strict compact closed version of $G\mathfrak{G}$. More precisely, the functor $F^{\mathrm{st}}: G\mathfrak{G} \to 1\operatorname{Cob}_{\mathfrak{G}}$ obtained by the above universal property of $G\mathfrak{G}$ is defined as follows. It maps every object X of $G\mathfrak{G}$ to the sequence of signs corresponding to the signed set P(X) (see [20, Section 3]). On arrows, it is defined just by replacing the source and the target by the corresponding sequences of signs. Namely, an arrow of $G\mathfrak{G}$ (see Ref. [20, Section 3]) is represented by a triple, which is essentially contained in the notion of \mathfrak{G} -cobordism. Hence, F^{st} maps an arrow (neglecting its source and target) to itself. It is straightforward to see that we have the following.

Proposition 5 The functor F^{st} : $G\mathfrak{G} \to 1\mathbf{Cob}_{\mathfrak{G}}$ is faithful.

In another words, to pass from $1\mathbf{Cob}_{\mathfrak{G}}$ to $G\mathfrak{G}$, one has to "decorate" the objects of $1\mathbf{Cob}_{\mathfrak{G}}$ with propositional formulae built in the language including single propositional letter, constant *I*, unary connective * and binary connective \otimes . However, this just disguises strict compact closed nature of $1\mathbf{Cob}_{\mathfrak{G}}$, which is intrinsic to this category.

Let $1\mathbf{Cob}_{\mathfrak{G}}^+$ be the category with the same objects as $1\mathbf{Cob}_{\mathfrak{G}}$, while the arrows of $1\mathbf{Cob}_{\mathfrak{G}}^+$ from *a* to *b* are the formal sums of arrows of $1\mathbf{Cob}_{\mathfrak{G}}$ from *a* to *b*. These formal sums may be represented by finite (possibly empty) multisets of \mathfrak{G} -cobordisms from *a* to *b*. Formally, a multiset of elements of a set *X* is a function from *X* to the set of natural numbers (including zero). Less formally, it is a set in which elements may have multiple occurrences.

We abuse the notation using the set brackets $\{, \}$ for multisets and by denoting a singleton multiset $\{f\}$ by f. Note that in this notation $\bigcirc + \bigcirc$, i.e., $\{\bigcirc, \bigcirc\}$ is not equal to $\bigcirc \bigcirc$, i.e., $\{\bigcirc\bigcirc\}$, where \bigcirc is a circular component with arbitrary label.

The identity arrow $\mathbf{1}_a : a \to a$ is the singleton multiset $\mathbf{1}_a : a \to a$, while the composition of $\{f_j : a \to b \mid j \in J\}$ and $\{f_k : b \to c \mid k \in K\}$ is

$$\{f_k \circ f_j \colon a \to c \mid j \in J, k \in K\}.$$

Again, because of too many roles of \emptyset in this paper, we denote the empty multiset of \mathfrak{G} -cobordisms from *a* to *b* by $0_{a,b}$, and call it *zero-arrow*. The existence of zero-arrows implies that every hom-set in 1**Cob**⁺_{\mathfrak{G}} is inhabited.

The category $1\mathbf{Cob}_{\mathfrak{G}}^+$ is enriched over the category **Cmd**. The addition in Hom (a, b) is the operation + (disjoint union) on multisets and the neutral is $0_{a,b}$.

Let $1\mathbf{Cob}_{\mathfrak{G}}^{\oplus}$ be the biproduct completion of $1\mathbf{Cob}_{\mathfrak{G}}^{+}$ constructed as follows (see Ref. [31, Section 5.1]). The objects of $1\mathbf{Cob}_{\mathfrak{G}}^{\oplus}$ are finite (possibly empty) sequences $(a_0, \ldots, a_{n-1}), n \ge 0$, of objects a_0, \ldots, a_{n-1} of $1\mathbf{Cob}_{\mathfrak{G}}$. (We abuse the notation by denoting a singleton sequence (a_0) by a_0 .) For example, (++-+-, o, +, --+, o) is an object of $1\mathbf{Cob}_{\mathfrak{G}}^{\oplus}$. (Here, according to our convention, o denotes the empty sequence of oriented points.)

The empty sequence of objects of $1\mathbf{Cob}_{\mathfrak{G}}$ plays the role of zero-object in $1\mathbf{Cob}_{\mathfrak{G}}^{\oplus}$, and for the above reasons, we denote it by **0** and not by \emptyset . Note the distinction between this object and the object presented by the singleton sequence *o* whose only member is the empty sequence of oriented points.

The arrows of $1\mathbf{Cob}_{\mathfrak{G}}^{\oplus}$ from (a_0, \ldots, a_{n-1}) to (b_0, \ldots, b_{m-1}) are the $m \times n$ matrices whose ij-entry is an arrow of $1\mathbf{Cob}_{\mathfrak{G}}^+$ from a_j to b_i . If m = 0, i.e., $(b_0, \ldots, b_{m-1}) = \mathbf{0}$, then the empty matrix is the unique arrow from $a = (a_0, \ldots, a_{n-1})$ to $\mathbf{0}$, and we denote it by $0_{a,0}$. We proceed analogously when n = 0.

The identity arrow $\mathbf{1}_a$ on $a = (a_0, \ldots, a_{n-1})$ in $1\mathbf{Cob}_{\mathfrak{G}}^{\oplus}$ is the $n \times n$ matrix with corresponding identity arrows of $1\mathbf{Cob}_{\mathfrak{G}}^+$ in the main diagonal and corresponding zero-arrows of $1\mathbf{Cob}_{\mathfrak{G}}^+$ elsewhere. The arrows are composed by the rule of matrix multiplication, save that the addition and multiplication in a field are replaced by addition in hom-sets and composition in the category $1\mathbf{Cob}_{\mathfrak{G}}^+$. For $a = (a_0, \ldots, a_{n-1})$ and $b = (b_0, \ldots, b_{m-1})$, we denote by $0_{a,b}$, or simply $0_{m \times n}$, the $m \times n$ matrix whose ij-entry is the zero-arrow $0_{a_j,b_i}$ of $1\mathbf{Cob}_{\mathfrak{G}}^+$. In the limit cases, when we compose the empty matrices $0_{a,0}$ and $0_{0,b}$, we define the result as the zero-matrix $0_{a,b}$.

Proposition 6 The category $1\mathbf{Cob}_{\mathfrak{G}}^{\oplus}$ has the structure of strict compact closed category with biproducts. The group of automorphisms of the object + in this category is isomorphic to \mathfrak{G} . Moreover, \dagger is definable in $1\mathbf{Cob}_{\mathfrak{G}}^{\oplus}$, which makes it dagger strict compact closed category with dagger biproducts, while the automorphisms of + are unitary.

Proof We define the compact closed structure on $1\mathbf{Cob}_{\mathfrak{G}}^{\oplus}$ as follows. The tensor product of objects (a_0, \ldots, a_{n-1}) and (b_0, \ldots, b_{m-1}) is the object $(a_0 \otimes b_0, \ldots, a_0 \otimes b_{m-1}, \ldots, a_{n-1} \otimes b_{m-1})$ of $1\mathbf{Cob}_{\mathfrak{G}}^{\oplus}$. If either n = 0 or m = 0, the result is zero-object **0**. The unit object is o. The tensor product of arrows of $1\mathbf{Cob}_{\mathfrak{G}}^{\oplus}$ is defined as the Kronecker product of matrices over a field, save that this time the multiplication in the field is replaced by the tensor product in the category $1\mathbf{Cob}_{\mathfrak{G}}^{\oplus}$.

The arrows α and λ are identities. For $a = (a_0, \ldots, a_{n-1})$ and $b = (b_0, \ldots, b_{m-1})$, the $(n \cdot m) \times (m \cdot n)$ matrix $\sigma_{a,b}$ (an arrow of $1\mathbf{Cob}_{\mathfrak{G}}^{\oplus}$) is defined as the permutation matrix representing the isomorphism between $V \otimes W$ and $W \otimes V$ for V being *n*-dimensional and W being *m* dimensional vector space, save that instead of the entries 1, we have arrows σ from $1\mathbf{Cob}_{\mathfrak{G}}^+$, with corresponding indices, and instead of entries 0, we have zero-arrows (i.e., empty multisets) of $1\mathbf{Cob}_{\mathfrak{G}}^+$ with corresponding indices. For example, if $a = (a_0, a_1, a_2)$ and $b = (b_0, b_1)$, the matrix $\sigma_{a,b}$ (with indices of zero-arrows omitted) is

1	σ_{a_0,b_0}	$0_{a_0\otimes b_1,b_0\otimes a_0}$	0	0	0	0)	
	0	0	σ_{a_1,b_0}	0	0	0	
	0	0	0	-	σ_{a_2,b_0}	0	
	0	σ_{a_0,b_1}	0	0	0	0	•
	0	0	0	σ_{a_1,b_1}	0	0	
	0	0	0	0	~	σ_{a_2,b_1}	

The operation * on objects of $1\mathbf{Cob}_{\mathfrak{G}}^{\oplus}$ is defined componentwise. The arrow η_a for $a = (a_0, \ldots, a_{n-1})$ is the $n^2 \times 1$ matrix with the singleton multiset η_{a_k} in the $k \cdot (n + 1)$ -th row, for $0 \le k < n$, and zero-arrows of $1\mathbf{Cob}_{\mathfrak{G}}^+$, with corresponding indices, elsewhere. The arrow ε_a is the $1 \times n^2$ matrix having ε_{a_k} in the $k \cdot (n + 1)$ -th column, for $0 \le k < n$, and zero-arrows of $1\mathbf{Cob}_{\mathfrak{G}}^+$, with corresponding indices, elsewhere. One can verify that the equalities A.1–A.12 hold in $1\mathbf{Cob}_{\mathfrak{G}}^{\oplus}$. Moreover, the arrows $u_{a,b}$, v and w_a defined in Sect. 2 are identities. Hence, $1\mathbf{Cob}_{\mathfrak{G}}^{\oplus}$ is a strict compact closed category.

The operation + on arrows from (a_0, \ldots, a_{n-1}) to (b_0, \ldots, b_{m-1}) is defined componentwise and zero-matrices are the neutrals for this operation. Equations A.10–A.12 hold, which guarantees that 1**Cob**^{\oplus}_{\mathfrak{S}} is enriched over **Cmd**.

For objects $a = (a_0, \ldots, a_{n-1})$ and $b = (b_0, \ldots, b_{m-1})$, the object $a \oplus b$ is the sequence

$$(a_0,\ldots,a_{n-1},b_0,\ldots,b_{m-1}).$$

The object **0** is the zero-object of $1\mathbf{Cob}_{\mathfrak{G}}^{\oplus}$ and it is the neutral for \oplus . For arrows $A_{m \times n}$ and $B_{p \times q}$ of $1\mathbf{Cob}_{\mathfrak{G}}^{\oplus}$, its *direct sum* $A \oplus B$ is the $(m + p) \times (n + q)$ matrix

$$\begin{pmatrix} A & 0_{m \times q} \\ 0_{p \times n} & B \end{pmatrix}.$$

For $a = (a_0, ..., a_{n-1})$ and $b = (b_0, ..., b_{m-1})$, the arrows $\pi^1_{a,b}, \pi^2_{a,b}, \iota^1_{a,b}$ and $\iota^2_{a,b}$ are defined as

$$\pi_{a,b}^{1} = \left(\mathbf{1}_{a} \ 0_{n \times m}\right), \quad \pi_{a,b}^{2} = \left(\mathbf{0}_{m \times n} \ \mathbf{1}_{b}\right),$$
$$\iota_{a,b}^{1} = \left(\mathbf{1}_{a} \\ 0_{m \times n}\right), \quad \iota_{a,b}^{2} = \left(\mathbf{0}_{n \times m} \\ \mathbf{1}_{b}\right).$$

After checking that the equalities A.13–A.19 hold in $1\mathbf{Cob}_{\mathfrak{G}}^{\oplus}$, one concludes that this category is strict compact closed with biproducts.

That the group of automorphisms of the object + in $1\mathbf{Cob}_{\mathfrak{G}}^{\oplus}$ is isomorphic to \mathfrak{G} is shown as follows. Every arrow from + to itself is a 1×1 matrix whose entry is a multiset of arrows of $1\mathbf{Cob}_{\mathfrak{G}}$ from the singleton sequence of oriented points + to itself. This multiset is a singleton in the case of an isomorphism, which follows from the fact that the composition in $1\mathbf{Cob}_{\mathfrak{G}}^{+}$ of a multiset of cardinality *n* with a multiset of cardinality *m* is a multiset of cardinality $n \cdot m$, and an isomorphism must be canceled to $1_{+}: + \rightarrow +$, which is the singleton multiset 1_{+} . Hence, every isomorphism from + to + in $1\mathbf{Cob}_{\mathfrak{G}}^{\oplus}$ is of the form $\psi: + \rightarrow +$, for ψ an arrow of $1\mathbf{Cob}_{\mathfrak{G}}$. Moreover, ψ must be an isomorphism in $1\mathbf{Cob}_{\mathfrak{G}}$. An arrow of $1\mathbf{Cob}_{\mathfrak{G}}$ from + to + consists of a single \mathfrak{G} -segment and several (possibly zero) \mathfrak{G} -circles. Since ψ is an isomorphism and \mathfrak{G} -circles are not cancelable, there are no \mathfrak{G} -circles in ψ and it could be identified with the underlying \mathfrak{G} -segment. The label of this segment is the element of \mathfrak{G} corresponding to the initial isomorphism of $1\mathbf{Cob}_{\mathfrak{G}}^{\oplus}$. It is evident that this correspondence is a one-to-one homomorphism.

The operation \dagger on arrows of $1\mathbf{Cob}_{\mathfrak{G}}^+$ is defined as

$$\left(\{f_j: a \to b \mid j \in J\}\right)^{\dagger} = \{f_j^{\dagger}: b \to a \mid j \in J\}.$$

The operation \dagger on a matrix representing an arrow of $1\mathbf{Cob}_{\mathfrak{G}}^{\oplus}$ is defined by transposing this matrix, and by applying the operation \dagger , defined above, to each of its entries. To verify that $1\mathbf{Cob}_{\mathfrak{G}}^{\oplus}$ is dagger strict compact closed with dagger biproducts, it remains to check that the equalities A.20–A.24 of Appendix A hold. The definition of \dagger in $1\mathbf{Cob}_{\mathfrak{G}}$ guarantees that the automorphisms of + are unitary.

Remark 2 For our purposes, it is useful to have a direct presentation of $\lceil f \rceil$, $\lfloor f \rfloor$, f_* and $\langle f_1, \ldots, f_n \rangle$ at least for arrows f, f_1, \ldots, f_n of 1**Cob**_G. The first three operations are defined as in 1**Cob** (the labels of \mathfrak{G} -components remain the same in the first two cases, while in the case of f_* the labels become the inverses of the initial labels). The last operation (see the definition of biproducts in Sect. 2) produces the $n \times 1$ matrix

$$\begin{pmatrix} f_1 \\ \vdots \\ f_n \end{pmatrix}.$$

Remark 3 By relying on the equalities 2.3, it is not difficult to show that the left distributivity isomorphism $v_{a,b,c}: (a \oplus b) \otimes c \to (a \otimes c) \oplus (b \otimes c)$ is the identity in the category $1\mathbf{Cob}_{\mathfrak{G}}^{\oplus}$. Similarly, if we assume that a is a singleton sequence, then by relying on the equalities 2.2, we can show that the right distributivity isomorphism $\tau_{a,b,c}: a \otimes (b \oplus c) \to (a \otimes b) \oplus (a \otimes c)$ is also the identity in $1\mathbf{Cob}_{\mathfrak{G}}^{\oplus}$.

Remark 4 By the universal property of the category \mathcal{F} from Sect. 4, there exists a unique functor $H: \mathcal{F} \to 1\mathbf{Cob}_{\mathfrak{G}}^{\oplus}$ that strictly preserves the compact closed structure with biproducts, for which Hp = + and for every $\gamma \in \Gamma$, $H\gamma$

is the \mathfrak{G} -cobordism from + to + given by one \mathfrak{G} -segment labeled by γ . The isomorphism of \mathcal{F} and \mathcal{F}^{\dagger} from the proof of Proposition 1 enables one to consider H as a functor from \mathcal{F}^{\dagger} to 1**Cob**^{\oplus} that strictly preserves the dagger compact closed structure with dagger biproducts.

Proposition 7 The functor $H: \mathcal{F} \to 1\mathbf{Cob}_{\mathfrak{G}}^{\oplus}$ is faithful.

Proof Let $f, g: a \to b$ be two arrows of \mathcal{F} such that Hf = Hg, and let $I_a = (\iota_a^0, \ldots, \iota_a^{n-1})$ and $\Pi_b = (\pi_b^0, \ldots, \pi_b^{m-1})$. By Corollary 1 and properties of biproducts, it suffices to show that, for every $0 \le i < m$ and $0 \le j < n$,

$$\pi_b^i \circ f \circ \iota_a^j = \pi_b^i \circ g \circ \iota_a^j. \tag{7.1}$$

By Corollary 3, both sides of 7.1 are expressible free of \oplus , ι , and π . By relying on the equalities A.11, A.12, A.27 and A.28, both sides are expressible as sums of terms, which are all free of \oplus , + and 0, ι , π -arrows. Here, the empty sum is denoted by $0_{a^j,b^i}$.

If one side of the above equality is equal to $0_{a^j,b^i}$, then it is mapped by *H* to the empty multiset. By functorial properties of *H*, the sum at the other side must be mapped by *H* to the empty multiset too, which means that this sum is empty, i.e., it is $0_{a^j,b^i}$.

It remains the case when for $n, m \ge 1$, the left-hand side of 7.1 is equal to $\sum_{k=1}^{n} f_k$ and the right-hand side of this equality is equal to $\sum_{k=1}^{m} g_k$, for f_k, g_k free of \oplus , + and 0, ι, π -arrows. We have that

$$\{Hf_k \mid 1 \le k \le n\} = H \sum_{k=1}^n f_k = H \sum_{k=1}^m g_k = \{Hg_k \mid 1 \le k \le m\},\$$

which means that n = m, and modulo some permutation of elements of these multisets, for every $1 \le k \le n$, $Hf_k = Hg_k$. The terms f_k and g_k belong entirely to the compact closed fragment generated by \mathfrak{G} . Hence, these terms represent arrows of a compact closed category $F\mathfrak{G}$ freely generated by \mathfrak{G} (see Ref. [20, Section 4]). The functor H restricts to $F\mathfrak{G}$ as the composition of an isomorphism (from $F\mathcal{A}$ to $G\mathcal{A}$, for \mathcal{A} being \mathfrak{G} ; see Ref. [20, Section 4]) and the faithful functor F^{st} of Proposition 5, which means that this restriction is faithful. We conclude that f_k and g_k represent the same arrow of $F\mathfrak{G}$, and hence of \mathcal{F} .

8 Validity of categorical quantum protocols

There are a lot of different frameworks for studying quantum protocols and checking its validity. Let us mention here only the recent studies Refs. [9] and [26] as the illustrative examples. To determine the requirements of quantum network protocols, the authors in Ref. [9] developed a software for the modeling and simulation of quantum networks (a network simulator for quantum information). In Ref. [26], the author showed (by estimating the validity of the quantum key distribution protocol as a paradigmatic situation) that an ensemble of states that can be distinguished by local operations and classical communication is more efficient for quantum information processing protocol than those states that cannot be discriminated deterministically in this way. Compared to these studies, our approach is more theoretical, and it relies on the work of Abramsky and Coecke [1].

It was suggested in Ref. [1] that compact closed categories with biproducts provide a generalization of von Neumann's presentation of quantum mechanics in terms of Hilbert spaces, [37]. Such an approach is called *categorical quantum mechanics*. For a survey of theory of categorical quantum mechanics, we recommend Ref. [2,35] and references therein.

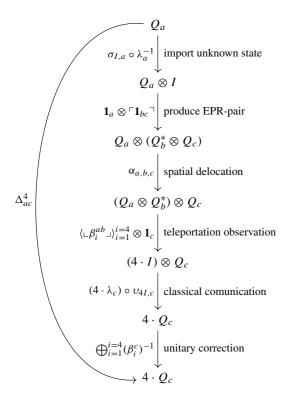
In this section, we use Proposition 7 to establish commutativity of diagrams in the category \mathcal{F} , which provides a verification of the corresponding protocols from the realm of categorical quantum mechanics. All the protocols verified in Ref. [1] require a compact closed category with dagger biproducts, possessing some additional structure. For the first two protocols below, this extra structure consists of an object Q (the *qubit*), an arrow from $4 \cdot I$ to $Q^* \otimes Q$ and a scalar *s* satisfying some conditions listed in [1, Section 9]. (Here we abbreviate $((I \oplus I) \oplus I) \oplus I$) by $4 \cdot I$, and more generally, $n \cdot a$ and $n \cdot f$ abbreviate the *n*-fold biproducts, associated to the left, of an object *a* and an arrow *f* respectively.)

However, the only additional structure upon a compact closed structure with dagger biproducts important for the verification diagrams consists of four unitary isomorphisms $\beta_1, \beta_2, \beta_3, \beta_4 \colon Q \to Q$. Hence, to establish that the verification diagrams are commutative in an arbitrary such category, i.e., that the categorical quantum protocols are correct, it suffices to establish their commutativity in the compact closed category \mathcal{F} with biproducts freely generated by the free group \mathfrak{G} on four generators. (Since β_1 is standardly taken to be identity, a group with three generators suffices.) The role of the generator p for objects of \mathcal{F} (see Sect. 4) belongs now to the qubit Q.

Our Proposition 7 enables one to check the commutativity of diagrams in \mathcal{F} by "drawing pictures" and this is the style of verification given below. We are aware that our graphical language is specific to the proposed framework, and it may have limited explanatory power for complex quantum phenomena. The qubit Q is interpreted in $1\mathbf{Cob}_{\mathfrak{G}}^{\oplus}$ as +. At some points, we have to draw matrices of pictures and this is done in the first example below, otherwise just the ij element of such a matrix is described. Of course, we focus on a certain class of protocols, as our goal is to show that there are cases where validity of quantum protocols can be checked using diagrams that are *geometric*, meaning that they correspond to manifolds (and a group element, but see the next section on ideas how to remove those elements).

8.1 Quantum teleportation

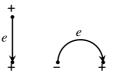
Quantum teleportation is a well-known quantum protocol [1,27]. Assume that Alice has a qubit in some state $|\psi\rangle$ and wishes to send this state to Bob without any knowledge of what this state is. This is done by taking an entangled pair of qubits (EPR-pair, $|\beta_1\rangle$) and sending one to Alice and another to Bob. Then Alice measures (in the Bell basis) her qubit and the qubit that is entangled with the one Bob has. In the next step, she communicates the result of the measurement to Bob, who applies unitary corrections to his qubit, depending on Alice's outcome. The final result is that Bob's qubit is in the same state as Alice's qubit originally was (Alice does not have a qubit in state $|\psi\rangle$ after this protocol is done).



(8.1)

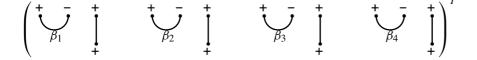
The correctness of the quantum teleportation protocol is expressed by commutativity of the diagram given in Ref. [1, Theorem 9.1]. One can easily factor the scalars out of both legs in this diagram just by appealing to the equalities 2.4–2.6. This makes the commutativity of Diagram 8.1 sufficient for the correctness of the protocol. We follow the terminology and notation introduced in Ref. [1] in this diagram.

Note that we treat Q_a , Q_b , and Q_c as three instances of the same object Q of \mathcal{F} . Also, $\Delta_{a,c}^4$ is an abbreviation for $\langle \mathbf{1}_Q, \mathbf{1}_Q, \mathbf{1}_Q, \mathbf{1}_Q, \mathbf{1}_Q \rangle$. For example, producing the EPR-pair means to apply the arrow $\mathbf{1}_Q \otimes \lceil \mathbf{1}_Q \rceil$, which is interpreted in 1**Cob**^{\oplus}_{\mathfrak{B}} as:

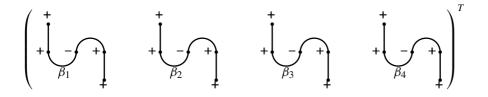


This is the first nontrivial step in the diagram 8.1. (Note that since $1\mathbf{Cob}_{\mathfrak{G}}^{\oplus}$ is a strict compact closed category, the steps called "import unknown state" and "spatial delocation" are interpreted as identities in this category.)

In drawings of \mathfrak{G} -cobordisms, when we interpret the arrows of the diagram 8.1 and the diagrams below, the orientation and the label *e* (denoting the neutral of \mathfrak{G}) will be omitted. As we noted at the beginning of this section, our group \mathfrak{G} is generated by the set $\Gamma = \{\beta_1, \beta_2, \beta_3, \beta_4\}$. The second nontrivial step in the diagram 8.1 is the teleportation observation, given by $\langle \square \beta_i \square \rangle_{i=1}^{i=4} \otimes \mathbf{1}_Q$, or in terms of arrows of $1\mathbf{Cob}_{\mathfrak{G}}^{\oplus}$:

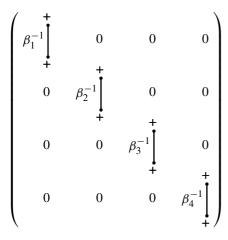


By composing this 4×1 matrix with the 1×1 matrix representing production of EPR-pair, we get

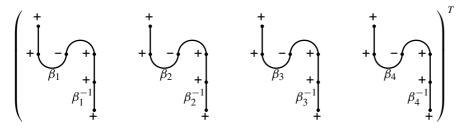


At this point, Alice had performed her measurement, and communicated the result to Bob using classical interchange of bits. By Remark 3, we know that the distributivity isomorphism $v_{4I,c}$ is the identity in 1Cob[⊕]_{\mathcal{C}}. This, together with the strictness of this category, makes the step named "classical communication" trivial, i.e. it is interpreted as identity.

Next, Bob applies unitary corrections, given by $\bigoplus_{i=1}^{i=4} (\beta_i)^{-1}$. In our matrix representation, \bigoplus corresponds to the direct sum of matrices. We, therefore, have the unitary correction



By composing the last two matrices, we get the final result



By stretching the diagrams, the group elements cancel out, and we are left with the diagonal $\Delta^4 = \langle \mathbf{1}_Q, \mathbf{1}_Q, \mathbf{1}_Q, \mathbf{1}_Q, \mathbf{1}_Q \rangle$.

8.2 Entanglement swapping

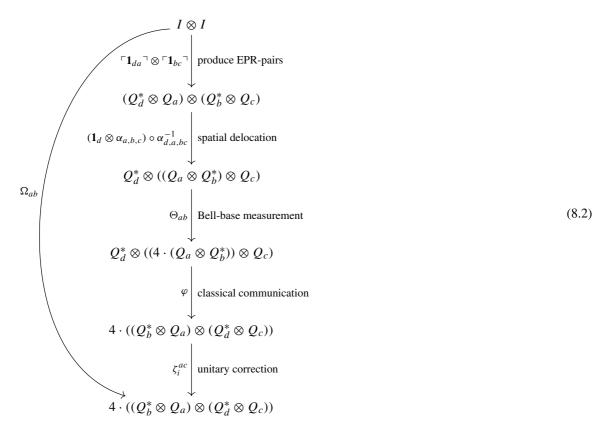
The idea of this protocol is to, starting with two pairs of mutually entangled qubits in EPR-states, obtain again two pairs of entangled states, but with different pairing. Assume Alice, as well as Bob, share a single EPR-pair with a third person, named Charlie. Then Charlie performs a measurement on his qubits, and via classical communication transfers information on his outcomes to other parties, upon which a unitary correction is applied. The net result of this protocol is that Alice and Bob share an entangled EPR-pair, while Charlie is left with another EPR-pair. We, thus, say that the entanglement is swapped. A complete description of this protocol in terms of categorical quantum mechanics is presented in Ref. [1, Theorem 9.3]. Again, as in Sect. 8.1, by relying on the equalities 2.4–2.6, one may completely neglect the role of scalars and just check the commutativity of the diagram 8.2 below for the correctness of this protocol.

Let $\tau: Q_d^* \otimes (4 \cdot ((Q_a \otimes Q_b^*) \otimes Q_c)) \to 4 \cdot (Q_d^* \otimes ((Q_a \otimes Q_b^*) \otimes Q_c))$ and $\upsilon: (4 \cdot ((Q_a \otimes Q_b^*)) \otimes Q_c \to 4 \cdot ((Q_a \otimes Q_b^*) \otimes Q_c))$ be distributivity isomorphisms, and let

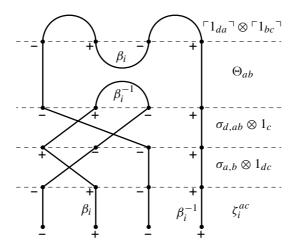
$$\begin{split} \gamma_{i} &= (\beta_{i})_{*}, \\ \mathbf{P}_{i} &= \lceil \gamma_{i} \rceil \circ \llcorner \beta_{i} \lrcorner, \\ \zeta_{i}^{ac} &= \bigoplus_{i=1}^{i=4} ((\mathbf{1}_{b}^{*} \otimes \beta_{i}) \otimes (\mathbf{1}_{d}^{*} \otimes \beta_{i}^{-1})), \\ \Theta_{ab} &= \mathbf{1}_{d}^{*} \otimes (\langle \mathbf{P}_{i} \rangle_{i=1}^{i=4} \otimes \mathbf{1}_{c}), \\ \varphi &= (4 \cdot ((\sigma_{ab} \otimes \mathbf{1}_{dc}) \circ \alpha_{ab,d,c}^{-1} \circ (\sigma_{d,ab} \otimes \mathbf{1}_{c}) \circ \alpha_{d,ab,c})) \circ \tau \circ (\mathbf{1}_{d} \otimes \upsilon), \end{split}$$

 $\Omega_{ab} = \langle \lceil \mathbf{1}_{ba} \rceil \circ \lceil \mathbf{1}_{dc} \rceil \rangle_{i=1}^{i=4}.$

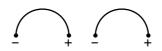
The commutativity of the diagram from Ref. [1, Theorem 9.3] justifies the correctness of the entanglement swapping protocol. By factoring the scalars out from the legs, it reduces to the following diagram.



The right-hand side of this diagram is represented in $1\mathbf{Cob}_{\mathfrak{G}}^{\oplus}$ by the 4×1 matrix whose *i*1-entry is the following \mathfrak{G} -cobordism (note that we ignore associativity and distributivity isomorphisms since they are identities).



By stretching the above diagram and canceling β_i and β_i^{-1} , we are left with the following \mathfrak{G} -cobordism.



On the other side, Ω_{ab} is represented in 1 **Cob**^{\oplus} by 4 × 1 matrix, whose *i*1-entry is exactly the above \mathfrak{G} -cobordism. Due to Proposition 7, this proves the commutativity of the diagram 8.2.

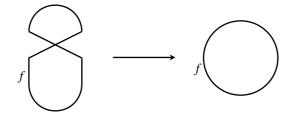
8.3 Superdense coding

In this section, we will apply our diagrammatic verification to another protocol, called superdense coding, [27] (sometimes referred to as dense coding). This quantum algorithm can be considered as the opposite of quantum teleportation. The idea is to transfer some amount of classical information using qubits. A review of this protocol can be found in Ref. [35], where its validity was shown in a similar manner.

The validity of this protocol is expressed in the categorical setting by the commutativity of a diagram in which some special scalars, namely traces of some arrows, occur. Every compact closed category can be lifted to the traced category by a suitable definition of a categorical trace. This can be achieved as follows. Let $f : a \to a$ be an arrow in a compact closed category. The scalar $Tr(f): I \to I$ is defined as¹

$$\operatorname{Tr}(f) = \varepsilon_a \circ (f \otimes a^*) \circ \sigma_{a^*,a} \circ \eta_a.$$
(8.3)

In terms of diagrams, we have



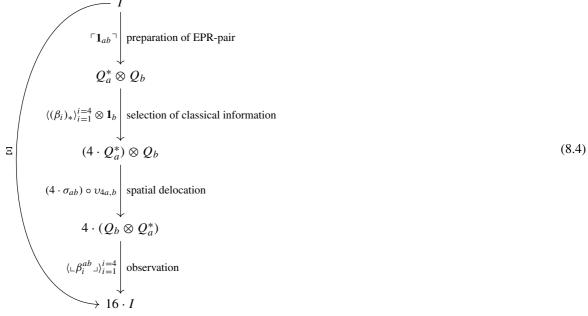
A category appropriate for the superdense coding requires the same structure as in the first two protocols. Moreover, the following conditions must be satisfied. If $i \neq j$, then $\text{Tr}(\beta_i \beta_j^{\dagger}) = 0_{I,I}$, and $\text{Tr}(\mathbf{1}_Q) \neq 0_{I,I}$ (see Appendix B for the details why we demand this condition to be satisfied). With this in mind, the arrow $\Xi: I \rightarrow 16 \cdot I$ defined as

$$(\operatorname{Tr}(\beta_1\beta_1^{\dagger}), \operatorname{Tr}(\beta_1\beta_2^{\dagger}), \operatorname{Tr}(\beta_1\beta_3^{\dagger}), \operatorname{Tr}(\beta_1\beta_4^{\dagger}), \dots, \operatorname{Tr}(\beta_4\beta_1^{\dagger}), \operatorname{Tr}(\beta_4\beta_2^{\dagger}), \operatorname{Tr}(\beta_4\beta_3^{\dagger}), \operatorname{Tr}(\beta_4\beta_4^{\dagger}))$$

is actually $\langle t, 0, 0, 0, 0, t, 0, 0, 0, t, 0, 0, 0, 0, t \rangle$, for $t = \text{Tr}(\mathbf{1}_Q)$. The assumption above also enables Bob to make a distinction between the four quadruples of scalars in this row.

¹ More generally, categorical trace corresponds to the partial trace in Hilbert-space picture, though we will not review this here, as our interest lies only in pure states.

Our task is to show that the following diagram, which verifies the superdense coding protocol, commutes.



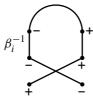
Here, again, the first step is the EPR-pair production, achieved by a cap diagram.



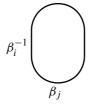
One qubit is located at Alice's point, and another at Bob's. Alice then applies an unitary transformation to her qubit, depending on the classical infromation she wants to communicate. This is achieved by $\langle (\beta_i)_* \rangle_{i=1}^{i=4} \otimes \mathbf{1}_b$. By composing the first two arrows, we get a 4×1 matrix, whose *i*1-entry is given by a following arrow.



Spatial delocation is represented by a transposition, and after applying it, we obtain a matrix with *i*1-entry given by



Finally, Alice sends her qubit to Bob, who performs an entangled state measurement, given by a suitable coname. The result is a 16×1 matrix, whose (4(i - 1) + j)1- element is given by the \mathfrak{G} -circle

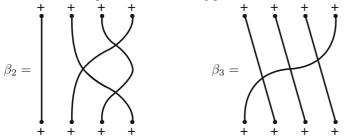


and the same matrix is obtained by interpreting the arrow $\Xi: I \to 16 \cdot I$ in the category $1\mathbf{Cob}_{\mathfrak{G}}^{\oplus}$. The additional assumptions on the compact closed structure, listed in the paragraph where the arrow Ξ is defined, enable Bob to distinguish between different Alice's messages.

9 Omitting labels

This section, in an informal way, illustrates possibility of elimination of labels assigned to cobordisms in graphical categories, introduced in Sect. 7, that serve as models for syntactical categories, introduced in Sect. 4. In the case when we have some additional equations concerning the unitary endomorphisms from Γ , it may produce a finite group of automorphisms of p. Then it is not necessary to increase the dimension of cobordisms, since such a group appears as a subgroup of a symmetric group S_n , for some n. One can represent p by n points and interpret every element of Γ as the corresponding permutation.

For example, the teleportation protocol requires the dihedral group D_4 , which is a subgroup of S_4 . This means that it is sufficient to define the syntactical category \mathcal{F}^{\dagger} so that Γ is a set of two elements that satisfy, besides the equalities 4.1 and 4.2, the equalities of the standard presentation of D_4 . In this case, the category 1**Cob**_{\mathfrak{G}}^{\oplus} should be replaced by a thickened version. This means that before labeling, every segment is replaced by four parallel threads (the diagram is thickened), and every label is replaced by the corresponding permutation of four threads. The two elements of Γ correspond to the following permutations.



However, there is no possibility to interpret an infinite group in such a way.

By increasing the dimension by one, according to the remark given in the penultimate paragraph on page 60 (after Proposition 1.4.9) of [22], the situation remains the same. This remark says that the only invertible 2-cobordisms are the permutation cobordisms.

Hence, for the interpretation of an infinite group generated by Γ , one has to consider 3-cobordisms. We rely here on Ref. [18, Definition 2.3] to introduce cobordisms that replace \mathfrak{G} -cobordisms from Sect. 7. Namely, for every orientation preserving homeomorphism $h: \Sigma_g \to \Sigma_g$, where Σ_g is a closed oriented surface of genus g, there is a cobordism ($\Sigma_g \times I$, f_0 , f_1), where $f_0(x) = (x, 0)$ and $f_1(x) = (h(x), 1)$. Two such cobordisms, corresponding to homeomorphisms h and h', respectively, are equivalent if and only if h and h' are pseudo-isotopic. According to Ref. [11], this is equivalent to the fact that h and h' are isotopic.

By applying technique introduced in Ref. [23], the cobordism ($\Sigma_g \times I$, f_0 , f_1) is equivalent to the cobordism (M, f_0 , f'_1), where M is $\Sigma_g \times I$ with some extra surgery, and $f'_1(x) = (x, 1)$. Here we will illustrate just the case of the group freely generated by one generator. In the case of a group freely generated by more generators, according to comments from the preceding paragraph, the results concerning free subgroups of the mapping class groups of surfaces, obtained in Ref. [4, 16, 17] are relevant.

In our example, we suggest to replace the \mathfrak{G} -segment labeled by the generator of \mathfrak{G} , i.e., a \mathfrak{G} -cobordism introduced in Sect. 7, by a 3-cobordism *C* obtained as follows. For T^2 being the 2-dimensional torus, the underlying manifold of *C* is $T^2 \times I$ with some additional surgery. Moreover, for $i \in \{0, 1\}$, the embeddings f_i are of the form $f_i(x) = (x, i)$. To present such a cobordism, we use the diagrammatical language introduced in Ref. [28] (see Fig. 1 for a presentation of *C*).

Roughly speaking, tubular neighborhoods of the red and the blue circle are removed from S^3 and the surgery along the black unknot is performed. Note that the twist of this unknot indicates the framing 1 of this surgery component.

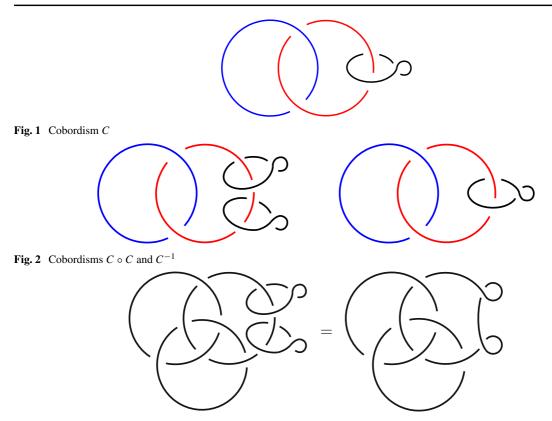


Fig. 3 Cobordism corresponding to the labeled \mathfrak{G} -circle

We refer to Ref. [28] for details of the interpretation of such diagrams. The rules for composing diagrams say that $C \circ C$ is presented by the diagram at the left-hand side of Fig. 2 and that the \mathfrak{G} -segment labeled by the inverse of the generator of \mathfrak{G} should be replaced by the cobordism C^{-1} illustrated at the right-hand side of Fig. 2.

The same rules say that the \mathfrak{G} -circle labeled by the product of the generator with itself should be replaced by the cobordism (closed 3-manifold) illustrated in Fig. 3.

Though this switching to dimension three could be less practical at some points, it could bring some new insight to the subject through the variant of Kirby calculus introduced in Ref. [13]. Our plan for a future work is to investigate this 3-dimensional calculus.

10 Concluding remarks

After the introduction of categorical quantum mechanics, it is natural to seek for a different dagger compact closed categories with biproducts, to check whether they can sustain quantum protocols, as quantum teleportation. The possible complication is the existence of a base. Abstractly, base can be defined using biproducts: we demand existence of unitary arrows $I \oplus I \rightarrow Q$. The problem with the category of cobordisms is that there does not seem to be enough options to construct the desired unitary morphism. This was alluded, in a slightly different context, in Ref. [5]. Luckily, to verify protocols as quantum teleportation, it is not mandatory to use the described morphism.

Furthermore, in low dimensions, it is hopeless to try to accommodate different unitary transformations present in quantum protocols as different cobordisms. To heal this problem, we introduced a group structure \mathfrak{G} . We believe that the approach suggested in the preceding section could provide a solution for these problems.

Of course, this raises some conceptual questions. First, by identifying the qubit state space with + (or –), we are not able to use our graphical language to define states, i.e., morphisms of the form $I \rightarrow Q$. In all mentioned quantum protocols, this was not an issue, as we used names to create entangled states, and this can be seen in 1**Cob**_{\mathfrak{G}}

language. To circumvent this issue, one could increase the dimension of cobordisms as suggested in the preceding section, or to take zero-dimensional spheres, i.e., the two element sequences + – to represent state spaces. Then one has the possibility to introduce morphisms that define states. Also, we can use this new type of qubits to define measurements on a single qubit, not just on an entangled pair. Considerations of this type could be of interest when dealing with single-particle protocols [10].

In addition to those limitations we already discussed, we stress that one drawback to our work is that it is not applicable to any quantum protocol. It would probably be hopeless to try dealing with complicated quantum systems whose entanglement structure is some complicated tensor network [6,36]. However, we still showed that there are quantum protocols, relying on the notion of entanglement, that are captured in this framework. On the other side, our work shows the possibility of connecting quantum mechanics and geometry, and this is the most important theoretical implication of the proposed framework. Furthermore, our work demonstrates that this relation can be made using the categorical approach to quantum mechanics and, therefore, hints that it could be fruitful to extend this approach to deal with other quantum systems, including quantum gravity.

Finally, we comment on the future directions of our work, continuing the discussion from the last section. We believe that it would be useful to obtain a three-dimensional structure that could distinguish between some different non-homeomorphic manifolds and that this structure could be used to check the validity of certain quantum protocols. In this light, our ambition is to add new cobordisms to our category that are not present in one-dimensional cases (as objects in this case are too simple). This way, we should be able to truly obtain quantum/geometry duality, which is, in the present case, spoiled by the introduction of group elements. Hopefully, pursuing this direction will tell us something new about the quantum world or at least emphasize some aspects that could be overlooked otherwise. One example is the notion of measurement in quantum mechanics. Even though this approach does not suggest a new way of looking at the measurement postulate of quantum mechanics, we notice that in our graphical language with matrices of pictures, the fact that the size of a matrix is enlarged after measurement resembles the many-world interpretation of quantum mechanics [12].

We conclude this section with a comment concerning the generality of quantum protocols brought by replacing the Hilbert spaces by objects of a compact closed category. It is known that (with minor provisos) all the onedimensional topological quantum field theories, i.e., functors from the category 1**Cob** to the category of finite dimensional vector spaces over a field, are faithful according to Ref. [32]. However, this does not mean that the whole 1**Cob**^{\oplus} could be faithfully represented by matrices over a field. On the other hand, since protocols do not use the full strength of 1**Cob**^{\oplus}, one could expect that some could be verified by relying on the matrix calculus (working again in the skeleton of **fdHilb** with chosen bases). This could be an advantage concerning computational issues of the problem.

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Data Availability This manuscript has no associated data.

Declarations

Conflict of interest The authors have no relevant financial or non-financial interests to disclose.

Ethical approval Not applicable.

Informed consent Not applicable.

(A.7)

Appendix

Appendix A The language and the equations for dagger compact closed categories with dagger biproducts

Our choice of a language for dagger compact closed categories with dagger biproducts is the one in which enrichment over **Cmd** is primitive and not derived from the biproduct structure. Such a language is situable for the proofs of our results. A dagger compact closed category with dagger biproducts \mathcal{A} consists of a set of objects and a set of arrows. There are two functions (*source* and *target*) from the set of arrows to the set of objects of \mathcal{A} . For every object *a* of \mathcal{A} there is the identity arrow $\mathbf{1}_a: a \to a$. The set of objects includes two distinguished objects *I* and 0. Arrows $f: a \to b$ and $g: b \to c$ compose to give $g \circ f: a \to c$, and arrows $f_1, f_2: a \to b$ add to give $f_1 + f_2: a \to b$. For every object *a* of \mathcal{A} , there is the object a^* , and for every pair of objects *a* and *b* of \mathcal{A} , there are the objects $a \otimes b$ and $a \oplus b$. Also, for every arrow $f: a \to b$, there is the arrow $f^{\dagger}: b \to a$, and for every pair of arrows $f: a \to a'$ and $g: b \to b'$ there are the arrows $f \otimes g: a \otimes b \to a' \otimes b'$ and $f \oplus g: a \oplus b \to a' \oplus b'$. In \mathcal{A} we have the following families of arrows indexed by its objects.

 $\begin{array}{ll} \alpha_{a,b,c} \colon a \otimes (b \otimes c) \to (a \otimes b) \otimes c, & \alpha_{a,b,c}^{-1} \colon (a \otimes b) \otimes c \to a \otimes (b \otimes c), \\ \lambda_a \colon I \otimes a \to a, & \lambda_a^{-1} \colon a \to I \otimes a, \\ \sigma_{a,b} \colon a \otimes b \to b \otimes a, & & \\ \eta_a \colon I \to a^* \otimes a, & \varepsilon_a \colon a \otimes a^* \to I \\ \pi_{a,b}^1 \colon a \oplus b \to a, & \iota_{a,b}^1 \colon a \to a \oplus b, \\ \pi_{a,b}^2 \colon a \oplus b \to b, & \iota_{a,b}^2 \colon b \to a \oplus b, \\ 0_{a,b} \colon a \to b. & & \\ \end{array}$

The arrows of A should satisfy the following equalities:

$$f \circ \mathbf{1}_a = f = \mathbf{1}_{a'} \circ f, \quad (h \circ g) \circ f = h \circ (g \circ f), \tag{A.1}$$

$$\mathbf{1}_a \otimes \mathbf{1}_b = \mathbf{1}_{a \otimes b}, \quad (f_2 \otimes g_2) \circ (f_1 \otimes g_1) = (f_2 \circ f_1) \otimes (g_2 \circ g_1), \tag{A.2}$$

$$((f \otimes g) \otimes h) \circ \alpha_{a,b,c} = \alpha_{a',b',c'} \circ (f \otimes (g \otimes h)),$$

$$\alpha_{a,b,c}^{-1} \circ \alpha_{a,b,c} = \mathbf{1}_{a \otimes (b \otimes c)}, \quad \alpha_{a,b,c} \circ \alpha_{a,b,c}^{-1} = \mathbf{1}_{(a \otimes b) \otimes c},$$
(A.3)

$$f \circ \lambda_a = \lambda_{a'} \circ (I \otimes f), \quad \lambda_a^{-1} \circ \lambda_a = \mathbf{1}_{I \otimes a}, \quad \lambda_a \circ \lambda_a^{-1} = \mathbf{1}_a, \tag{A.4}$$

$$(g \otimes f) \circ \sigma_{a,b} = \sigma_{a',b'} \circ (f \otimes g), \quad \sigma_{b,a} \circ \sigma_{a,b} = \mathbf{1}_{a \otimes b}, \tag{A.5}$$

$$\alpha_{a\otimes b,c,d} \circ \alpha_{a,b,c\otimes d} = (\alpha_{a,b,c} \otimes d) \circ \alpha_{a,b\otimes c,d} \circ (a \otimes \alpha_{b,c,d}), \tag{A.6}$$

$$\lambda_{a\otimes b} = (\lambda_a \otimes b) \circ \alpha_{I,a,b},$$

$$\alpha_{c,a,b} \circ \sigma_{a \otimes b,c} \circ \alpha_{a,b,c} = (\sigma_{a,c} \otimes b) \circ \alpha_{a,c,b} \circ (a \otimes \sigma_{b,c}), \tag{A.8}$$

$$(a^* \otimes \varepsilon) \circ \alpha_{a^*, a, a^*}^{-1} \circ (\eta \otimes a^*) = \sigma_{I, a^*}, \quad (\varepsilon \otimes a) \circ \alpha_{a, a^*, a} \circ (a \otimes \eta) = \sigma_{a, I},$$
(A.9)

$$f_1 + (f_2 + f_3) = (f_1 + f_2) + f_3, \quad f_1 + f_2 = f_2 + f_1, \quad f + 0_{a,a'} = f,$$
 (A.10)

$$(g_1 + g_2) \circ f = g_1 \circ f + g_2 \circ f, \quad g \circ (f_1 + f_2) = g \circ f_1 + g \circ f_2, \tag{A.11}$$

$$0_{a',b} \circ f = 0_{a,b}, \quad f \circ 0_{b,a} = 0_{b,a'}, \tag{A.12}$$

$$\mathbf{1}_a \oplus \mathbf{1}_b = \mathbf{1}_{a \oplus b}, \quad (f_2 \oplus g_2) \circ (f_1 \oplus g_1) = (f_2 \circ f_1) \oplus (g_2 \circ g_1), \tag{A.13}$$

$$(f \oplus g) \circ \iota^1_{a,b} = \iota^1_{a',b'} \circ f, \quad (f \oplus g) \circ \iota^2_{a,b} = \iota^2_{a',b'} \circ g, \tag{A.14}$$

$$f \circ \pi^{1}_{a,b} = \pi^{1}_{a',b'} \circ (f \oplus g), \quad g \circ \pi^{2}_{a,b} = \pi^{2}_{a',b'} \circ (f \oplus g), \tag{A.15}$$

$$\pi_{a,b}^{1} \circ \iota_{a,b}^{1} = \mathbf{1}_{a}, \quad \pi_{a,b}^{2} \circ \iota_{a,b}^{2} = \mathbf{1}_{b},$$
(A.16)

$$\pi_{a,b}^2 \circ \iota_{a,b}^1 = 0_{a,b}, \quad \pi_{a,b}^1 \circ \iota_{a,b}^2 = 0_{b,a}, \tag{A.17}$$

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$$\iota_{a\,b}^{1} \circ \pi_{a\,b}^{1} + \iota_{a\,b}^{2} \circ \pi_{a\,b}^{2} = \mathbf{1}_{a \oplus b}. \tag{A.18}$$

$$0_{0,0} = \mathbf{1}_0,$$
 (A.19)

$$\mathbf{1}_{a}^{\dagger} = \mathbf{1}_{a}, \quad (g \circ f)^{\dagger} = f^{\dagger} \circ g^{\dagger}, \quad f^{\dagger \dagger} = f, \tag{A.20}$$

$$f \otimes g)^{\dagger} = f^{\dagger} \otimes g^{\dagger}, \tag{A.21}$$

$$\alpha_{a,b,c}^{\dagger} = \alpha_{a,b,c}^{-1}, \quad \lambda_a^{\dagger} = \lambda_a^{-1}, \quad \sigma_{a,b}^{\dagger} = \sigma_{b,a}, \tag{A.22}$$

$$\varepsilon^{\dagger} = \sigma_{a^*,a} \circ \eta, \tag{A.23}$$

$$(\pi_{a,b}^{1})' = \iota_{a,b}^{1}, \quad (\pi_{a,b}^{2})' = \iota_{a,b}^{2}.$$
(A.24)

The following equalities are derivable from A.1–A.24:

$$(f \oplus g)' = f' \oplus g', \tag{A.25}$$

$$(f+g)^{\dagger} = f^{\dagger} + g^{\dagger}, \quad 0^{\dagger}_{a,b} = 0_{b,a}$$
 (A.26)

$$f \otimes (g_1 + g_2) = (f \otimes g_1) + (f \otimes g_2), \quad (f_1 + f_2) \otimes g = (f_1 \otimes g) + (f_2 \otimes g), \tag{A.27}$$

$$f \otimes 0_{b,b'} = 0_{a \otimes b,a' \otimes b'} = 0_{a,a'} \otimes g.$$
(A.28)

Appendix B Scalars and probability amplitudes

As firmly laid, quantum mechanics is based on complex vector spaces (Hilbert spaces, to be more precise). Implied in this structure is the notion of scalars, that correspond here to the field of complex numbers. In categorical language, one can define scalars more abstractly [1,20]. A *scalar* is a morphism $s : I \rightarrow I$. It can be proved that the hom-set Hom (I, I), for a compact closed category, is a commutative monoid, therefore justifying further this structure's name.

In 1**Cob**, the scalars correspond to closed, one-dimensional manifolds, and the only candidate for such a structure is a finite collection of circles S^1 (as denoted on the left-hand side of the following picture). In 1**Cob**_{\mathfrak{G}}, we have \mathfrak{G} -circles; topological circles dressed with group elements (right-hand side of the following picture). Due to the compact closed structure of this category, there is a natural interpretation of those circles. Namely, any compact closed category can be lifted to a traced category by a suitable definition of a categorical trace (see Sect. 8.3 for the definition).



That closed loops should be connected with traces is not limited to a categorical approach to quantum mechanics. Even when considering Feynman diagrams in quantum electrodynamics, fermions loops are accompanied by a trace in spinorial indices. Moreover, in TQFT, we are customed to the fact that closing manifold by gluing the outward future to inward past (if possible), results in a trace, that for a cylinder, i.e. the identity, simply gives the dimension of the respective Hilbert space.

Furthermore, as explained in [7], these traces correspond to the probability weights of different branches. This is further confirmed by a Hilbert-space picture computations. Recall that one reason we have scalars (different from the multiplicative unit) is normalization on states. In order to get the probabilistic interpretation, according to the Born rule, we must insist on normalized states. For a state $\beta_{00}|0\rangle \otimes |0\rangle + \beta_{01}|0\rangle \otimes |1\rangle + \beta_{10}|1\rangle \otimes |0\rangle + \beta_{11}|1\rangle \otimes |1\rangle$, we have that its norm squared is given by $|\beta_{00}|^2 + |\beta_{01}|^2 + |\beta_{10}|^2 + |\beta_{11}|^2 = \text{Tr}(\beta^{\dagger}\beta)$, where β is a 2 × 2 matrix whose components are β_{ij} constants. Therefore, we conclude that traces are as important in this set-up as in traditional Hilbert-space formulation. When dealing with quantum protocols, one usually takes β to be proportional to Pauli sigma matrices. (Extended) Pauli matrices are defined as

$$\sigma_0 = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}, \quad \sigma_1 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \quad \sigma_2 = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}, \quad \sigma_3 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

We see that those matrices are unitary, self-adjoint and satisfy $\text{Tr}(\sigma_i \sigma_j) = 2\delta_{ij}$, where δ_{ij} is a Kronecker delta symbol (equal to one if i = j and zero otherwise). In order to make the connection with the Bell basis, introduced in Sect. 1, we take $\beta_1 = \sigma_0$, $\beta_2 = \sigma_1$, $\beta_3 = \sigma_3$ and $\beta_4 = -i\sigma_2$. This implies that we have $\text{Tr}(\beta_i \beta_j^{\dagger}) = 2\delta_{ij}$, with the usual definition of matrix adjoint.

However, in order to check whether two diagrams commute, it is usually straightforward to include scalars into consideration. One can then just neglect this issue of scalars and work without explicitly using them (as done previously). They are, of course, needed if one is to obtain probabilities for different outcomes of a measurement, but in our work (and related work of [1,7]) this is not a primary task.

References

- Abramsky, S., Coecke, B.: A categorical semantics of quantum protocols. In: Proceedings of the 19th Annual IEEE Symposium on Logic in Computer Science, LICS 2004, pp. 415–42. IEEE Computer Society Press5 (2004)
- Abramsky, S., Coecke, B.: Categorical quantum mechanics. In: Engesser, K., Gabbay, D.M., Lehmann, D. (eds.) Handbook of Quantum Logic and Quantum Structures, vol. 2, pp. 261–323. Elsevier, Oxford (2008)
- 3. Al-Raeei, M.: Applying fractional quantum mechanics to systems with electrical screening effects. Chaos Soliton Fract. **150**, 111209 (2021)
- 4. Anderson, J.W., Aramayona, J., Shackleton, K.J.: Free subgroups of surface mapping class groups. Conform. Geom. Dyn. 11, 229–321 (2007)
- 5. Baez, J.C.: Quantum quandaries: a category theoretic perspective. In: Rickles, D., French, S., Saatsi, J. (eds.) The Structural Foundations of Quantum Gravity, pp. 240–265. Oxford University Press, Oxford (2006)
- 6. Bridgeman, J.C., Chubb, C.T.: Hand-waving and interpretive dance: an introductory course on tensor networks. J. Phys. A Math. Theor. **50**, 223001 (2017)
- Coecke, B.: Kindergarten quantum mechanics: lecture notes. In: G. Adenier, A. Khrennikov and T.M. Nieuwenhuizen (eds.) Quantum Theory: Reconsiderations of the Foundations—3, AIP Conference Proceedings, vol. 810, pp. 81–98 (2005)
- 8. Coecke, B., Duncan, R.: Interacting quantum observables: categorical algebra and diagrammatics. New J. Phys. 13, 043016 (2011)
- 9. Coopmans, T., et al.: NetSquid, a NETwork Simulator for QUantum Information using Discrete events. Commun. Phys. 4, 164 (2021)
- 10. Del Santo, F., Dakić, B.: Two-way communication with a single quantum particle. Phys. Rev. Lett. 120, 060503 (2018)
- 11. Epstein, D.B.A.: Curves on 2-manifolds and isotopies. Acta Math. 115, 83–107 (1966)
- 12. Everett, H.: Relative state. Formulation of quantum mechanics. Rev. Mod. Phys. 29, 454–462 (1957)
- 13. Femić, B., Grujić, V., Obradović, J., Petrić, Z.: A calculus for S³-diagrams of manifolds with boundary, available at arXiv (2022)
- Harlow, D.: TASI lectures on the emergence of bulk physics in AdS/CFT. In: Proceedings of Theoretical Advanced Study Institute Summer School 2017, Boulder, Colorado, Proceedings of Science (2018)
- 15. Heunen, C.: Categorical Quantum Models and Logics. Pallas Publications, Amsterdam University Press, Amsterdam (2009)
- Ishida, A.: The structure of subgroup of mapping class groups generated by two Dehn twists. Proc. Jpn. Acad. Ser. A Math. Sci. 72, 240–241 (1996)
- Ivanov, N.V.: Subgroups of Teichmüller modular groups. American Mathematical Society, Translations of Mathematical Monographs, vol. 115 (1992)
- 18. Juhász, A.: Defining and classifying TQFTs via surgery. Quant. Topol. 9, 229-321 (2018)
- 19. Kauffman, L.: Jr. Lomonaco, Topological quantum information theory. In: Proceedings of Symposia in Applied Mathematics, p. 68 (2012)
- 20. Kelly, G.M., Laplaza, M.L.: Coherence for compact closed categories. J. Pure Appl. Algebra 19, 193-213 (1980)
- 21. Kirby, R.: A calculus for framed links in S^3 . Invent. Math. 45, 35–56 (1978)
- 22. Kock, J.: Frobenius Algebras and 2D Topological Quantum Field Theories. Cambridge University Press, Cambridge (2003)
- 23. Lickorish, W.B.R.: A representation of orientable combinatorial 3-manifolds. Ann. Math. 76, 531–540 (1962)
- 24. Maldacena, J.: The Large N limit of superconformal field theories and supergravity. Adv. Theor. Math. Phys. 2, 231–252 (1998)
- 25. Maldacena, J., Susskind, L.: Cool horizons for entangled black holes. Fortschr. Phys. 61, 781–811 (2013)
- 26. Nandi, S.: Catalysing quantum information processing task using LOCC distinguishability. Pramana J. Phys. **95**, 155 (2021)
- Nielsen, M., Chuang, I.: Quantum Computation and Quantum Information: 10th, Anniversary. Cambridge University Press, Cambridge (2010)
- 28. Nikolić, J., Petrić, Z., Zekić, M.: A diagrammatic presentation of the category 3Cob. RM 79, 165 (2024)

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- 29. Petrić, Z., Zekić, M.: Coherence for closed categories with biproducts. J. Pure Appl. Algebra 225, 106533 (2021)
- 30. Ryu, S., Takayanagi, T.: Aspects of holographic entanglement entropy. J. High Energy Phys. 2006(08), 045 (2006)
- Selinger, P.: Dagger compact closed categories and completely positive maps. In: Quantum Programming Languages, Electronic Notes in Theoretical Computer Science, vol. 170, pp. 139-163. Elsevier (2007)
- 32. Telebaković Onić, S.: On the faithfulness of 1-dimensional topological quantum field theories. Glasnik Mat. 55, 67-83 (2020)
- 33. Tong, D., Wong, K.: Monopoles and Wilson Lines. J. High Energy Phys. 2014, 48 (2014)
- 34. Turaev, V.G.: Quantum Invariants of Knots and 3-Manifolds. De Gruyter, Berlin/New York (2010)
- 35. Vicary, J.: Higher Quantum Theory, arXiv preprint (2012). arXiv:1207.4563
- 36. Vidal, G.: Entanglement renormalization. Phys. Rev. Lett. 99, 220405 (2007)
- 37. von Neumann, J.: Mathematische Grundlagen der Quantenmechanik. Springer, Berlin (1932). (Mathematical Foundations of Quantum Mechanics, Princeton University Press, English translation (1955))

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